EXISTENCE AND UNIQUENESS OF DIFFUSIONS ON THE JULIA SETS OF MISIUREWICZ-SIERPINSKI MAPS

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ABSTRACT. We study the balanced resistance forms on the Julia sets of Misiurewicz-Sierpinski maps, which are self-similar resistance forms with equal weights. In particular, we use a theorem of Sabot to prove the existence and uniqueness of balanced forms on these Julia sets. We also provide an explorative study on the resistance forms on the Julia sets of rational maps with periodic critical points.

1. INTRODUCTION

The study of diffusion processes on fractals emerged as an independent research field in the late 80's. Initial interest in such processes came from mathematical physicists working in the theory of disordered media [1, 19, 36]. On self-similar sets, the pioneering works are the constructions of Brownian motions on the Sierpinski gasket [15, 24] originated by Kusuoka and Goldstein independently and later [7] by Barlow and Perkins, and on the Sierpinski carpet [4] by Barlow and Bass. See [25] for an equivalent but different construction put forth by Kusuoka and Zhou at about the same time. The Sierpinski gasket is finitely ramified, meaning that the fractal can be disconnected by the removal of finitely many points, and the construction was later extended to wider families of fractals, such as the nested fractals [26] by Lindstrøm, and the post-critically finite (p.c.f.) self-similar sets [20, 21] by Kigami. Due to the porous structure of the fractals, the diffusion processes move slower on average than a standard Brownian motion on $\mathbb{R}^d$, see [5, 7, 13, 16, 23] for the associated transition density estimates. See books [3, 22, 38] for systematic explorations of the subject and more bibliographies.

On p.c.f. self-similar sets, Kigami [20, 21] showed that Dirichlet forms can be constructed as limits of electrical networks on approximating graphs. The construction relies on determining a proper form on the initial graph, whose existence and uniqueness in general is a difficult and fundamental problem in fractal analysis. On nested fractals, Lindstrøm proved that there always exists a symmetric diffusion process [26]. The problem was also investigated on nested fractals and p.c.f. self-similar sets by Metz [29] and Sabot [35] respectively. In particular, Sabot ingeniously proved the uniqueness of a symmetric diffusion process of equal weights on nested fractals by introducing the notion of *preserved relations*. See also [30] by Metz, and [31] by Peirone for short proofs. For the uniqueness problem on p.c.f. self-similar sets, see [17, 35] for some sufficient conditions, and see [32, 33] by Peirone for an effectively computable necessary and sufficient condition once a solution to (3.2) is known. The uniqueness theorem

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for non-p.c.f. self-similar sets is more difficult, see [6] for a positive answer for the generalized Sierpinski carpets.

In this paper, we utilize Sabot's techniques on nested fractals to study the existence and uniqueness of self-similar diffusions on a new class of finitely ramified fractals, the *Julia sets of Misiurewicz-Sierpinski maps*, as introduced in [12]. Let

$$R_{\lambda,n,m}(z) = z^n + \frac{\lambda}{z^m}, \quad n \geq 2, m \geq 1, \lambda \in \mathbb{C} \setminus \{0\}$$

be a rational map. A point $c \in \hat{\mathbb{C}} := \mathbb{C} \cup \{\infty\}$ is a *critical point* of $R_{\lambda,n,m}$, if $R_{\lambda,n,m}$ fails to be injective in any neighbourhood of c. In particular, the set of critical points of $R_{\lambda,n,m}$ consists of

$$C = \{c \in \mathbb{C} : R'_{\lambda,n,m}(c) = 0\},$$

∞ and 0 (if $m \geq 2$). We call $R_{\lambda,n,m}$ a *Misiurewicz-Sierpinski map* (*MS map* for short) if:

(MS1). each critical point in C is strictly preperiodic; and
(MS2). each critical point in C is on the boundary of the immediate attracting basin of ∞;

Here, in (MS1), we say a point z is strictly preperiodic if z is not periodic but eventually periodic under the iteration of $R_{\lambda,n,m}$, i.e. $R^k_{\lambda,n,m}(z) \neq z$, $\forall k \geq 1$, and $R^k_{\lambda,n,m}(z) = R^{k+l}_{\lambda,n,m}(z)$ for some $k, l \geq 1$. In (MS2), the attracting basin of ∞ is $\{z \in \hat{\mathbb{C}} : \lim_{k\to\infty} R^k_{\lambda,n,m}(z) = \infty\}$, and the immediate attracting basin of ∞ is the unique connected component of the attracting basin of ∞ containing ∞.

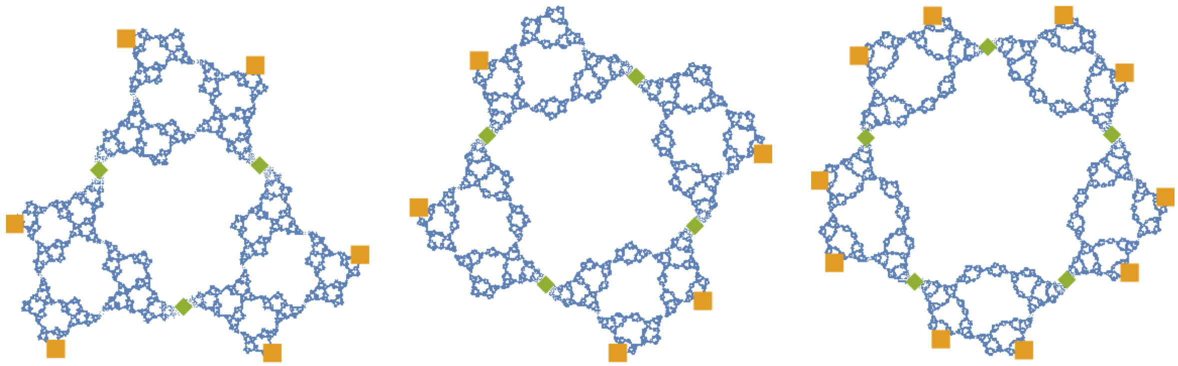
The existence of MS maps has been guaranteed in [34]. Indeed, for fixed m, n, there are infinitely many different (in the sense of homeomorphic) Julia sets $K_{\lambda,n,m}$ (K_λ for short, see Section 2 for the exact definition) associated with MS maps. See [12] and [34].

The dynamics and topological properties of the Julia sets K_λ associated with the MS maps were studied in [10, 12]. In general, K_λ is a generalized Sierpinski gasket, in the sense that it is a limit set obtained by a similar recursive process generating the Sierpinski gasket.

Recall that a Sierpinski gasket is obtained by starting with a triangle, dividing it into four small triangles and deleting the interior of the central one, then recursively repeating the same operation on the remaining triangles and continuing. A MS Julia set K_λ is obtained similarly but applied instead to the closed unit disk D as starting set and by removing homeomorphic copies of polygons of $N = m + n$ sides. First, we remove from D the interior of a homeomorphic N-polygon with only its corners lying in the boundary of D, and so we are left with a connected set composed by the union of N homeomorphic copies of D, denoted as $D_1, D_2, \cdots, D_N$. Next, we keep on removing the interior of a homeomorphic copy of N-polygon from each of D_i with only corners lying in the boundary of D_i. We repeat the process to the limit to obtain the Julia set K_λ.

Due to this construction, there is a natural p.c.f. self-similar structure on K_λ, with an iterated function system (i.f.s. for short) $\{F_i\}_{i=1}^{m+n}$, which will be described in Section 2. See Figure 1 for some examples of such fractals, where the green blocks denote critical points, and the orange blocks denote orbits of critical points under the iteration of $R_{\lambda,n,m}$.

The family of Julia sets of MS maps provides us a rich class of fractals. In particular, for fixed m, n and MS parameters λ, τ, if $\tau \notin \{\lambda e^{\frac{2k\pi i}{n-1}}, \bar{\lambda} e^{\frac{2k\pi i}{n-1}} : 1 \leq k \leq n-1\}$, K_λ and K_τ are not topologically equivalent [12]. The self-similar structure can be complicated depending



$m = 1, n = 2$ $\lambda \approx -0.0380422 + 0.42623i$ $m = 2, n = 2$ $\lambda \approx -0.0196286 + 0.275378i$ $m = 3, n = 2$ $\lambda \approx -0.01295 + 0.204151i$

FIGURE 1. Examples of MS Julia sets.

on the choice of λ. Despite these difficulties, we will prove the existence and uniqueness of balanced resistance forms on such Julia sets.

Theorem 1. *Let $R_{\lambda,n,m}$ be a MS map, and K_λ be the associated Julia set. There exists a unique resistance form $(\mathcal{E}_\lambda, \mathcal{F}_\lambda)$ on K_λ, such that*

$$\mathcal{E}_\lambda(f) = \eta \sum_{i=1}^{m+n} \mathcal{E}_\lambda(f \circ F_i), \quad \forall f \in \mathcal{F}_\lambda,$$

for some constant $\eta > 1$. We call such a form a balanced form.

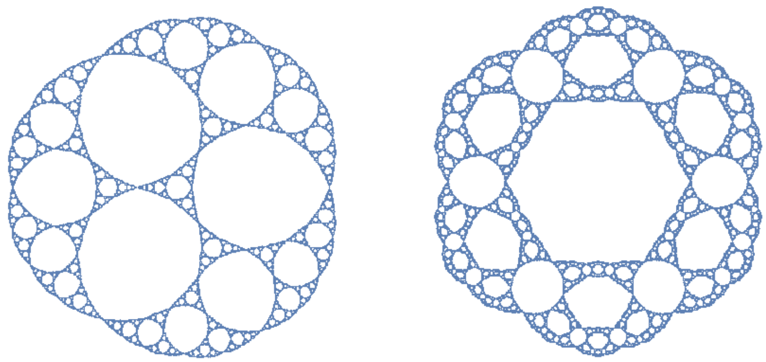
We remark that throughout the paper, a unique resistance form means unique up to a multiplicative constant. Though we are considering resistance forms with equal weights, our story has some essential differences with that of nested fractals.

1. The uniqueness is a little stronger than that on nested fractals in that symmetry of the form is not required by the problem. The same conclusion is not true in general on nested fractals, for example the Vicsek sets admit infinitely many different resistance forms with equal weights [28, 35].

2. Compared to the MS Julia sets, nested fractals have a larger symmetry group, big enough so that any pair of the boundary vertices are permuted by some element, and the existence can be proven with a fixed point argument [26]. On the other hand, our proof of existence will use the full strength of Sabot's techniques. In particular, the proof depends crucially on the dynamics of $R_{\lambda,n,m}$ on K_λ. One main difficulty is to find all possible non-trivial preserved $\mathcal{G}$-relations.

We mention that there have been several previous works studying the resistance forms on Julia sets of polynomial maps [2, 14, 37], but the methods and goals are quite different from those in this paper. Additionally, at the end of this work, we study other Julia sets associated with rational maps, specifically whose critical set possesses a real fixed point. See Figure 2 for an illustration. The resistance forms on such Julia sets admit graph-directed structures.

We briefly introduce the structure of the paper. In Section 2, we will introduce the p.c.f. self-similar structures and some dynamical properties of the Julia sets of MS maps. Section



$m = 1, n = 2$ $m = 3, n = 3$

FIGURE 2. Julia sets of $R_{\lambda,n,m}$ with a fixed critical point.

3 will be our main section, where we prove Theorem 1. This section will be divided into 4 parts. In the first part, we review the construction of resistance forms, and introduce Sabot's theorem. In the second part, we prove the existence of resistance forms. In the third part, we prove the uniqueness of resistance forms. Lastly, at the end of Section 3, we provide some examples. In Section 4, we present explorative results on the Julia sets of rational maps with a fixed critical point. Some rough discussions on the existence and non-existence of forms will be provided.

Throughout this paper, we will write R_λ instead of $R_{\lambda,n,m}$ if no confusion is likely.

2. A review of Misiurewicz-Sierpinski maps

In this section, we briefly review some simple properties of MS maps, and introduce self-similar structures on their associated Julia sets. Readers can find more details in [12].

Let R be a rational map. A periodic point z of R is called repelling if $|(R^k)'(z)| > 1$ for some $k \geq 1$. The *Julia set* of R is the closure of all repelling periodic points of R. Each connected component of the complement of the Julia set is called a *Fatou component* of R.

Let

$$R_{\lambda,n,m}(z) = z^n + \frac{\lambda}{z^m}, \quad n \geq 2, m \geq 1, \lambda \in \mathbb{C} \setminus \{0\}$$

be a MS map, and K_λ be the associated Julia set. Let $C = \{c \in \mathbb{C} : R'_\lambda(c) = 0\}$ be the set of critical points, excluding the poles at 0 and ∞, as in Section 1, i.e.

$$C = \big\{c \in \mathbb{C} : nc^{n-1} - m\frac{\lambda}{c^{m+1}} = 0\big\}.$$

We call C the *critical set* of R_λ. Clearly, we have

$$\#C = m + n \quad \text{ and } e^{\frac{2\pi i}{m+n}} C = C,$$

i.e. C admits $(m+n)$-fold rotational symmetry around $z = 0$. In the above formula, the letter i in the term $e^{\frac{2\pi i}{m+n}}$ stands for the imaginary unit. Denote β_λ the boundary of the immediate attracting basin of ∞, and τ_λ the closure of $R_\lambda^{-1}(\beta_\lambda) \setminus \beta_\lambda$, i.e. the boundary of the Fatou component containing 0. Note that β_λ and τ_λ are Jordan curves [10, 12].

Lemma 2.1. *(a).* $R_\lambda(e^{\frac{2\pi i}{m+n}}z) = e^{\frac{2n\pi i}{m+n}}R_\lambda(z)$, *for any* $z \in \mathbb{C}$.
(b). $\bigcup_{k=1}^{\infty} R_\lambda^k(C) \cap \bigcup_{k=0}^{\infty} R_\lambda^{-k}(C) = \emptyset$.
(c). $C \subset \beta_\lambda \cap \tau_\lambda \subset K_\lambda$.

Proof. (a) follows from a direct computation.

(b). Suppose there is a point $z \in \bigcup_{k=1}^{\infty} R_\lambda^k(C) \cap \bigcup_{k=0}^{\infty} R_\lambda^{-k}(C)$. It follows that there are $c_1, c_2 \in C$ and $k \geq 1$ such that $R_\lambda^k(c_1) = c_2$. Consider the sequence $\{R_\lambda^{lk}(c_1)\}_{l\geq 1}$. By (a) and the fact $e^{\frac{2\pi i}{m+n}}C = C$, we have $\{R_\lambda^{lk}(c_1)\}_{l\geq 1} \subset C$, which implies that there is a periodic point $c \in C$, a contradiction to (MS1). Thus (b) follows.

(c) is proved in [12] in the case $m = n = 2$, but obviously holds in general by an almost identical proof. We only give a sketch here. Let B_λ be the immediate attracting basin of ∞, and T_λ be the Fatou component of R_λ containing 0. Since $R_\lambda^{-1}(\infty) = \{0, \infty\}$, we have $R_\lambda^{-1}(B_\lambda) = B_\lambda \cup T_\lambda$, and β_λ, τ_λ are the boundaries of B_λ and T_λ, respectively. Since $C \subset \beta_\lambda$ by (MS2), we have $R_\lambda(C) \subset R_\lambda(\beta_\lambda) = \beta_\lambda$. Note that R_λ is locally two-to-one at each point in C. This implies that $C \subset \beta_\lambda \cap \tau_\lambda$. Noting the boundaries of Fatou components are contained in the Julia set, we have $\beta_\lambda \cap \tau_\lambda \subset K_\lambda$. □

We will need the good geometry of K_λ based on the conditions (MS1) and (MS2).

Proposition 2.2. *A MS Julia set* K_λ *is a connected compact set. The removal of* C *from* K_λ *disconnects* K_λ *into exactly* $m + n$ *components. In addition,* $e^{\frac{2\pi i}{m+n}}K_\lambda = K_\lambda$, *i.e.* K_λ *admits* $(m + n)$*-fold rotational symmetry around* $z = 0$.

Proof. The Julia set K_λ of rational maps of the form $R_\lambda(z) = z^n + \frac{\lambda}{z^m}$ (without assuming (MS1),(MS2)) has been considered in [11]. A main result is that K_λ is connected if and only if $R_\lambda^2(C)$ is not contained in the immediate attracting basin of ∞. Clearly, for a MS map R_λ, the condition (MS2) guarantees that $C \subset K_\lambda$ and it is not contained in the attracting basin of ∞. So the MS Julia set K_λ is connected. The second statement follows from Lemma 2.1 (c), see also Theorem 3.3 in [12]. The $(m + n)$-fold symmetry of K_λ (without assuming (MS1),(MS2)) follows from Lemma 1.3 in [11]. □

Due to Proposition 2.2, we denote the $m+n$ components of $K_\lambda \setminus C$ as $\{\mathring{K}_{\lambda,1}, \cdots, \mathring{K}_{\lambda,m+n}\}$. For each $i \in \{1, 2, \cdots, m + n\}$, we let $K_{\lambda,i}$ be the closure of $\mathring{K}_{\lambda,i}$ and call it a 1-*cell* of K_λ. By the dynamical property of R_λ (see [12, p. 15], for a detailed description), for each i, the map R_λ is a homeomorphism from $K_{\lambda,i}$ to K_λ. We denote F_i the i-th branch of R_λ^{-1} which is a homeomorphism from K_λ to $K_{\lambda,i}$. We thus have

$$K_{\lambda,i} = F_i K_\lambda, \quad K_\lambda = \bigcup_{i=1}^{m+n} F_i K_\lambda.$$

In addition, observing from Proposition 2.2 and using the symmetry of K_λ, we can see that C consists of exactly those points connecting $K_{\lambda,i}$'s.

Corollary 2.3. $C = \bigcup_{i\neq j} F_i K_\lambda \cap F_j K_\lambda$.

Proof. First, for each pair $i \neq j$, it is easy to see $F_i K_\lambda \cap F_j K_\lambda \subset C$. Otherwise, one would have $\mathring{K}_{\lambda,i} \cap \mathring{K}_{\lambda,j} = (F_i K_\lambda \setminus C) \cap (F_j K_\lambda \setminus C) \neq \emptyset$, a contradiction to the fact that $\{\mathring{K}_{\lambda,1}, \cdots, \mathring{K}_{\lambda,m+n}\}$ are connected components of $K_\lambda \setminus C$.

Second, we need to show that for each $c \in C$, there are $i \neq j$ such that $c \in F_i K_\lambda \cap F_j K_\lambda$. Clearly, by Proposition 2.2, there are $c' \in C$, $i' \neq j'$ such that $c' \in F_{i'} K_\lambda \cap F_{j'} K_\lambda$. By the rotational symmetry of C, we have $c = e^{\frac{2k\pi i}{m+n}} c'$ for some $k \in \mathbb{Z}$. In addition, $K_\lambda \setminus C$ has also the rotational symmetry, $e^{\frac{2k\pi i}{m+n}} (K_\lambda \setminus C) = K_\lambda \setminus C$, so there are i, j such that $i \neq j$, $\mathring{K}_{\lambda,i} = e^{\frac{2k\pi i}{m+n}} \mathring{K}_{\lambda,i'}$ and $\mathring{K}_{\lambda,j} = e^{\frac{2k\pi i}{m+n}} \mathring{K}_{\lambda,j'}$. This implies that $c \in F_i K_\lambda \cap F_j K_\lambda$ as desired. □

We will show that $\mathcal{L} := \big(K_\lambda, \{1, 2, \cdots, m+n\}, \{F_i\}_{i=1}^{m+n}\big)$ described above is a post-critically finite (p.c.f.) self-similar structure. This notion was introduced by Kigami in [21], see also [22]. To be precise, we recall the p.c.f. self-similar structure as follows. Let K be a compact metric space, S be a finite set, and $\{F_i\}_{i\in S}$ be a collection of continuous injections from K to itself. Then $\big(K, S, \{F_i\}_{i\in S}\big)$ is called a *self-similar structure* if there exists a continuous surjection $\Lambda : S^{\mathbb{N}} \to K$ such that $F_i \circ \Lambda = \Lambda \circ \sigma_i$, where for each $i \in S$, $\sigma_i : S^{\mathbb{N}} \to S^{\mathbb{N}}$ is defined by $\sigma_i(\omega_1\omega_2\cdots) = i\omega_1\omega_2\cdots$ for each $\omega_1\omega_2\cdots \in S^{\mathbb{N}}$. For a self-similar structure $\mathcal{L} = \big(K, S, \{F_i\}_{i\in S}\big)$, define $C_{\mathcal{L}} = \bigcup_{i,j\in S, i\neq j} F_i K \cap F_j K$, $\mathcal{C}_{\mathcal{L}} = \Lambda^{-1} C_{\mathcal{L}}$, and $\mathcal{P}_{\mathcal{L}} = \bigcup_{n\geq 1} \sigma^n \mathcal{C}_{\mathcal{L}}$, where $\sigma : S^{\mathbb{N}} \to S^{\mathbb{N}}$ is the shift map defined by $\sigma(\omega_1\omega_2\cdots) = \omega_2\omega_3\cdots$. Then $\mathcal{L}$ is called a *p.c.f. self-similar structure* if $\#\mathcal{P}_{\mathcal{L}} < \infty$ [21, 22]. The set $V_0(\mathcal{L}) := \Lambda(\mathcal{P}_{\mathcal{L}})$ is called the *boundary* of K.

Proposition 2.4. *$\mathcal{L} = \big(K_\lambda, \{1, 2, \cdots, m+n\}, \{F_i\}_{i=1}^{m+n}\big)$ is a p.c.f. self-similar structure.*

Proof. First, according to Proposition 3.6 in [12], we have

$$\operatorname{diam}(F_{\omega_1} F_{\omega_2} \cdots F_{\omega_k} K_\lambda) \to 0, \text{ as } k \to \infty,$$

for any infinite word $\omega \in \{1, 2, \cdots, m+n\}^{\mathbb{N}}$. This provides a continuous surjective coding map $\Lambda : \{1, 2, \cdots, m+n\}^{\mathbb{N}} \to K_\lambda$ defined by

$$\{\Lambda(\omega)\} = \bigcap_{k=1}^{\infty} F_{\omega_1} F_{\omega_2} \cdots F_{\omega_k} K_\lambda, \quad \forall \omega \in \{1, 2, \cdots, m+n\}^{\mathbb{N}}.$$

Obviously, $F_i \circ \Lambda = \Lambda \circ \sigma_i$ for each $i \in \{1, 2, \cdots, m+n\}$. Thus $\mathcal{L}$ is a self-similar structure.

Next, define $\mathcal{C} = \Lambda^{-1} C$, $\mathcal{P} = \bigcup_{n\geq 1} \sigma^n \mathcal{C}$ and $V_0 = \Lambda(\mathcal{P})$. It is easy to check that $V_0 = \bigcup_{k=1}^{\infty} \bigcup_{w\in\{1,2,\cdots,m+n\}^k} F_{w_1}^{-1} \cdots F_{w_k}^{-1} C$, and thus

$$V_0 = \bigcup_{k=1}^{\infty} R_\lambda^k(C) \quad \text{ and } \mathcal{P} = \Lambda^{-1}(V_0).$$

By (MS1), each point in C is strictly preperiodic, so that $\#V_0 < \infty$. Now we claim that

$$\#\mathcal{P} = \#V_0,$$

and thus $\#\mathcal{P} < \infty$ also. In fact, by Lemma 2.1 (b), we have $V_0 \cap \bigcup_{k=0}^{\infty} R_\lambda^{-k}(C) = \emptyset$. On the other hand,

$$\begin{aligned}\{z \in K_\lambda : \#\Lambda^{-1}(z) > 1\} &= \bigcup_{k=0}^{\infty} \bigcup_{w\in\{1,2,\cdots,m+n\}^k} \bigcup_{i\neq j} F_{w_1} \cdots F_{w_k} F_i K_\lambda \cap F_{w_1} \cdots F_{w_k} F_j K_\lambda \\ &= \bigcup_{k=0}^{\infty} \bigcup_{w\in\{1,2,\cdots,m+n\}^k} F_{w_1} \cdots F_{w_k}(C) = \bigcup_{k=0}^{\infty} R_\lambda^{-k}(C).\end{aligned}$$

This implies that $V_0 \subset \{z \in K_\lambda : \#\Lambda^{-1}(z) = 1\}$, which implies $\#\mathcal{P} = \#V_0 < \infty$.

□

From now on, we call $V_0 = \bigcup_{k=1}^{\infty} R_\lambda^k(C)$ the *boundary* of K_λ. See Figure 1 for some examples of K_λ, with C (green blocks) and V_0 (orange blocks) marked.

Remark. An important fact about the boundary V_0 of K_λ is that

$$F_i K_\lambda \cap F_j K_\lambda = F_i V_0 \cap F_j V_0, \quad \forall i \neq j. \tag{2.1}$$

This follows from a general property of p.c.f. self-similar sets [22]. We will frequently use this fact in later developments.

We list some more topological or dynamical facts about K_λ shown in [10, 12], which will play the key role in Section 3.

- *The 'ring' shape of K_λ.*

By suitably ordering $K_{\lambda,i} = F_i K_\lambda$ and critical points $c_i \in C$, we have the property that

$$\{c_{i-1}, c_i\} = F_i K_\lambda \cap C = F_i K_\lambda \cap \bigcup_{j \neq i} F_j K_\lambda, \quad \text{for } 1 \leq i \leq m+n, \tag{2.2}$$

and

$$c_i = e^{\frac{2\pi i}{m+n}} c_{i-1}, \qquad \text{for } 1 \leq i \leq m+n,$$

where we use cyclic notation $m+n=0$. Readers should be aware that here (also elsewhere) the letter i in the term $e^{\frac{2\pi i}{m+n}}$ always stands for the imaginary unit, not the index. The 1-cells of K_λ form a 'ring' shape, see Figure 3 for an illustration.

The 'ring' shape of K_λ provides more details in addition to Proposition 2.2. In particular, one can see that

$$F_i K_\lambda = e^{\frac{2\pi i}{m+n}} F_{i-1} K_\lambda, \text{ for } 1 \leq i \leq m+n, \tag{2.3}$$

with cyclic notation $m+n=0$. In fact, by the rotational symmetry of C and K_λ, one would have $e^{\frac{2\pi i}{m+n}} F_i K_\lambda = F_j K_\lambda$ for some j, and $j = i+1$ by (2.2).

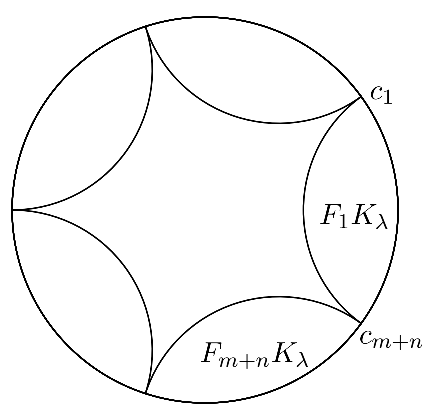


FIGURE 3. An illustration of level-1 cells.

- *Dynamics of R_λ on the outside boundary of K_λ.*

As before, let β_λ be the boundary of the immediate attracting basin of ∞. The critical set $C \subset \beta_\lambda$, disconnects β_λ into $m+n$ components, and each is contained in a cell $F_iK_\lambda, 1 \leq i \leq m+n$.

Furthermore, as shown in [12], $R_\lambda : \beta_\lambda \to \beta_\lambda$ is conjugate to a simple dynamic on the unit circle $\mathbb{T}$. More precisely, let $\mathbb{T}$ be the unit circle defined as

$$\mathbb{T} = \mathbb{R}/\mathbb{Z} = \{[r] = r + \mathbb{Z} : r \in \mathbb{R}\}.$$

For $n \geq 2$, we define $\Phi_n : \mathbb{T} \to \mathbb{T}$ by

$$\Phi_n([\theta]) = [n\theta].$$

Lemma 2.5 ([12]). *There is a homeomorphism $\psi_{\lambda,n,m} : \beta_\lambda \to \mathbb{T}$ (ψ_λ for short) such that*

$$\psi_\lambda \circ R_\lambda = \Phi_n \circ \psi_\lambda.$$

In addition, $\psi_\lambda(e^{\frac{2\pi i}{m+n}} z) = \psi_\lambda(z) + \frac{1}{m+n}$ by symmetry.

We denote by θ_λ the unique element of $\psi_\lambda(C)$ in $[0, \frac{1}{m+n})$. Since each $c \in C$ is strictly preperiodic, we have $\theta_\lambda \in \mathbb{Q}$. In addition,

$$\begin{cases} \psi_\lambda(C) = \{[\theta_\lambda + \frac{l}{m+n}] : 0 \leq l \leq m+n-1\}, \\ \psi_\lambda(V_0) = \{[n^k(\theta_\lambda + \frac{l}{m+n})] : k \geq 1, 0 \leq l \leq m+n-1\}. \end{cases} \tag{2.4}$$

Readers can read [12], Section 3 for more details, and a discussion on how θ_λ determines the dynamics.

Example 2.6. *(a). The first image in Figure 1 is the Julia set associated to $R_{\lambda,2,1}$ with $\lambda \approx -0.0380422 + 0.42623i$. For this simple example, we have $\theta_\lambda = \frac{1}{12}$. Thus,*

$$\psi_\lambda(C) = \{[\frac{1}{12}], [\frac{5}{12}], [\frac{3}{4}]\}, \text{ and } \psi_\lambda(V_0) = \{[0], [\frac{1}{6}], [\frac{1}{3}], [\frac{1}{2}], [\frac{2}{3}], [\frac{5}{6}]\}.$$

(b). The second image in Figure 1 is the Julia set associated to $R_{\lambda,2,2}$ with $\lambda \approx -0.0196286-0.275378i$, for which we have $\theta_\lambda = \frac{3}{16}$ and

$$\psi_\lambda(C) = \{[\frac{3}{16}], [\frac{7}{16}], [\frac{11}{16}], [\frac{15}{16}]\}, \text{ and } \psi_\lambda(V_0) = \{[0], [\frac{3}{8}], [\frac{7}{8}], [\frac{3}{4}], [\frac{1}{2}]\}.$$

In particular, this example shows that V_0 may not be rotational symmetric.

- *The topological structure of higher order cells of K_λ.*

It is often useful to consider smaller cells. Let's focus on a 1-cell F_iK_λ, which is bounded by $F_i\beta_\lambda$. Clearly F_iK_λ is disconnected into exactly $m+n$ components, if we remove F_iC. Another observation is that $\psi_\lambda(\beta_\lambda \cap F_iK_\lambda)$ is a closed arc of length $\frac{1}{m+n}$, which contains exactly n points in $\psi_\lambda(F_iC)$, see [12]. That means we have

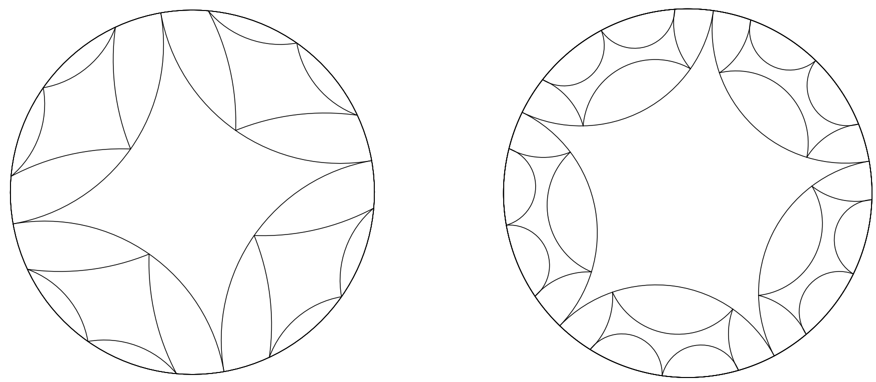
Lemma 2.7 ([12]). *For each $1 \leq i \leq m+n$,*

$$\#F_iC \cap \beta_\lambda = n, \quad \#F_iC \setminus \beta_\lambda = m.$$

With this in mind, we can sketch the level-2 cells, as illustrated in Figure 4.

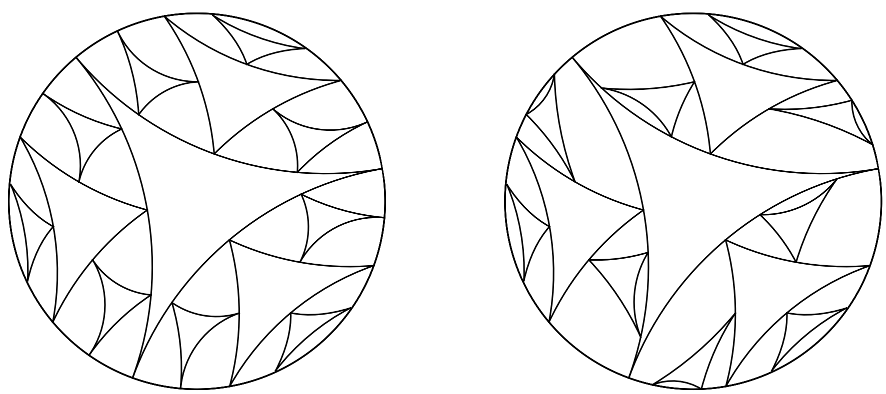
However, to accurately sketch the higher level cells, we need the exact value of θ_λ. There could exist several cases, see Figure 5. In general, the exact structure can depend in a complicated way on θ_λ.

We end this section by enumerating some results about MS parameters λ and θ_λ.



$m = 1, n = 3$ $m = 2, n = 3$

FIGURE 4. An illustration of the level-2 cells.



$m = 1, n = 2, \theta_\lambda = \frac{1}{6}$ $m = 1, n = 2, \theta_\lambda = \frac{1}{12}$

FIGURE 5. An illustration of the level-3 cells.

1). ([12, 39]) The set of parameter λ associated to MS maps is a dense subset of the boundary of the locus of connectedness of the family $\{R_{\lambda,n,m}\}$ (with fixed choice of n, m).
2). ([12]) For any two different MS maps $R_{\lambda,n,m}$ and $R_{\tau,n,m}$ with $\tau \notin \{\lambda e^{\frac{2k\pi i}{n-1}}, \bar{\lambda} e^{\frac{2k\pi i}{n-1}} : 1 \leq k \leq n-1\}$, their respective Julia sets are not topologically equivalent.
3). ([34]) For the $m = n$ case, there is a MS parameter λ for any strictly preperiodic θ_λ (under the dynamics of Φ_n).

3. Existence and uniqueness of balanced resistance forms

In this section, we consider the existence and uniqueness of self-similar diffusions on the Julia sets of Misiurewicz-Sierpinski maps. In [22], the concept of resistance forms is introduced, which in many cases describe local regular Dirichlet forms.

Let X be a set, and $l(X)$ be the space of all real-valued functions on X. A pair $(\mathcal{E}, \mathcal{F})$ is called a *(non-degenerate) resistance form* on X if it satisfies the following conditions:
(RF1). *$\mathcal{F}$ is a linear subspace of $l(X)$ containing constants and $\mathcal{E}$ is a nonnegative symmetric quadratic form on $\mathcal{F}$; $\mathcal{E}(u) := \mathcal{E}(u,u) = 0$ if and only if u is constant on X.*
(RF2). *Let $\sim$ be an equivalent relation on $\mathcal{F}$ defined by $u \sim v$ if and only if $u - v$ is constant on X. Then $(\mathcal{F}/\sim, \mathcal{E})$ is a Hilbert space.*
(RF3). *For any finite subset $V \subset X$ and for any $v \in l(V)$, there exists $u \in \mathcal{F}$ such that $u|_V = v$.*
(RF4). *For any $p, q \in X$, $r(p,q) := \sup\{\frac{|u(p)-u(q)|^2}{\mathcal{E}(u)} : u \in \mathcal{F}, \mathcal{E}(u) > 0\}$ is finite.*
(RF5). (Markov property) *If $u \in \mathcal{F}$, then $\bar{u} = \min\{\max\{u, 0\}, 1\} \in \mathcal{F}$ and $\mathcal{E}(\bar{u}) \leq \mathcal{E}(u)$.*

Since each cell of the Julia set is a copy of itself under compositions of homeomorphisms $\{F_i\}_{i=1}^{m+n}$, and the 1-cells $\{F_i K_\lambda\}_{i=1}^{m+n}$ are all the same size, a natural choice of resistance forms on K_λ are the *balanced resistance forms*, defined as follows.

Definition 3.1. *Let $R_{\lambda,n,m}$ be a MS map and let K_λ be the corresponding Julia set. We say a resistance form $(\mathcal{E}, \mathcal{F})$ on K_λ is balanced if there exists a positive constant η such that*

$$\mathcal{E}(f) = \eta \sum_{i=1}^{m+n} \mathcal{E}(f \circ F_i), \quad \forall f \in \mathcal{F}.$$

The main purpose of this paper is to prove the existence and uniqueness of a balanced resistance form on K_λ.

Theorem 3.2. *Let R_λ be a MS map, and K_λ be the corresponding Julia set. Then there is a unique balanced resistance form $(\mathcal{E}_\lambda, \mathcal{F}_\lambda)$ on K_λ. In addition, the unique balanced resistance form has rotational symmetry,*

$$\mathcal{E}_\lambda(f(e^{\frac{2\pi i}{m+n}} \bullet)) = \mathcal{E}_\lambda(f), \quad \forall f \in \mathcal{F}_\lambda.$$

In addition, by using well established results in [22], we can easily see the following result.

Theorem 3.3. *Let R_λ be a MS map, K_λ be the corresponding Julia set, and $(\mathcal{E}_\lambda, \mathcal{F}_\lambda)$ be the corresponding balanced resistance form. Let μ be a Radon measure on K_λ. Then $(\mathcal{E}_\lambda, \mathcal{F}_\lambda)$ becomes a local regular Dirichlet form on $L^2(K_\lambda, \mu)$.*

3.1. **The Theorem of Sabot.** We first demonstrate that the problem of finding a balanced resistance form (or more generally, a self-similar resistance form) can be transferred to a nonlinear fixed point problem on a finitely dimensional space.

Notation. Let X be a set equipped with a resistance form $(\mathcal{E}, \mathcal{F})$, and $V \subset X$ be a subset of finitely many points. We define the *restriction* of $(\mathcal{E}, \mathcal{F})$ to V by

$$\mathcal{E}|_V(f) = \inf\{\mathcal{E}(\tilde{f}) : \tilde{f}|_V = f, \tilde{f} \in \mathcal{F}\}, \quad \forall f \in l(V).$$

Note that $(\mathcal{E}|_V, l(V))$ is a resistance form on V by the polarization identity [22]. For $f \in l(V)$, we denote the unique (*harmonic*) *extension* of f with minimal energy by $H_{\mathcal{E},V} f$, so that

$$\mathcal{E}(H_{\mathcal{E},V} f) = \mathcal{E}|_V(f).$$

Though we do not highlight $\mathcal{F}$ in the above notations, the constructions depend on both $\mathcal{E}$ and $\mathcal{F}$.

Now, we assume that $(\mathcal{E},\mathcal{F})$ is a balanced resistance form on the Julia set K_λ. For each $f\in l(V_1)$, where

$$V_1 := \bigcup_{i=1}^{m+n} F_i(V_0),$$

we have

$$\mathcal{E}|_{V_1}(f)=\mathcal{E}(H_{\mathcal{E},V_1}f)=\eta\sum_{i=1}^{m+n}\mathcal{E}\big((H_{\mathcal{E},V_1}f)\circ F_i\big)\geq\eta\sum_{i=1}^{m+n}\mathcal{E}\big(H_{\mathcal{E},V_0}(f\circ F_i)\big)=\eta\sum_{i=1}^{m+n}\mathcal{E}|_{V_0}(f\circ F_i).$$

The other direction inequality also holds,

$$\mathcal{E}|_{V_1}(f)=\mathcal{E}(H_{\mathcal{E},V_1}f)\leq\mathcal{E}(g)=\eta\sum_{i=1}^{m+n}\mathcal{E}\big(H_{\mathcal{E},V_0}(f\circ F_i)\big)=\eta\sum_{i=1}^{m+n}\mathcal{E}|_{V_0}(f\circ F_i).$$

where g is the extension of f from $l(V_1)$ to $\mathcal{F}$ such that $g\circ F_i=H_{\mathcal{E},V_0}(f\circ F_i)$. Thus we have $(H_{\mathcal{E},V_1}f)\circ F_i=H_{\mathcal{E},V_0}(f\circ F_i)$, and

$$\mathcal{E}|_{V_1}(f)=\eta\sum_{i=1}^{m+n}\mathcal{E}|_{V_0}(f\circ F_i). \tag{3.1}$$

Notice that the equation $\big(\mathcal{E}|_{V_1}\big)|_{V_0}=\mathcal{E}|_{V_0}$ always holds, so that (3.1) becomes an identity of $\mathcal{E}|_{V_0}$.

On the other hand, if there exists a resistance form $\mathcal{D}$ on V_0 (we omit the domain $l(V_0)$ hereafter for simplicity) such that

$$T\mathcal{D}:=\mathcal{D}^{(1)}|_{V_0}=\eta^{-1}\mathcal{D}, \tag{3.2}$$

where $\mathcal{D}^{(1)}$ is the resistance form on V_1 defined by $\mathcal{D}^{(1)}(f)=\sum_{i=1}^{m+n}\mathcal{D}(f\circ F_i)$, and η is a positive constant, then there exists a unique balanced resistance form $(\mathcal{E},\mathcal{F})$ on K_λ such that $\mathcal{D}=\mathcal{E}|_{V_0}$.

Remark. For the purpose of using rotational symmetry later, we sometimes enlarge V_0 to $\tilde{V}_0:=\bigcup_{l=0}^{m+n-1}e^{\frac{2l\pi i}{m+n}}V_0$ and consider $\mathcal{D}$ on $\tilde{V}_0$ instead. Clearly, $\tilde{V}_1:=\bigcup_{i=1}^{m+n}F_i\tilde{V}_0=R_\lambda^{-1}\tilde{V}_0$ is also rotationally symmetric. We note that it follows from (2.4) $\tilde{V}_0=V_0$ when $m+n$ and n are coprime.

From equation (3.2), the problem amounts to a fixed point problem of the map T on the projective space of resistance forms (Dirichlet forms) on V_0 (or $\tilde{V}_0$). This problem is of fundamental importance in the study of diffusions on finitely ramified fractals. A famous and pioneering work on the existence of a solution was written by Lindstrøm for nested fractals [26], which are a class of highly symmetric p.c.f. fractals. Later, the problem of uniqueness for nested fractals was solved by Sabot in his celebrated work [35]. Moreover, Sabot raised a general theorem on both the existence (also non-existence) and the uniqueness of $\mathcal{D}$ to the solution of (3.2). Throughout the rest paper, without explicitly mention, a unique solution to (3.2) means unique up to a multiplicative constant.

We will utilize Sabot's theorem in our situation. Let's introduce some definitions from [35].

Consider a general p.c.f. self-similar structure $\mathcal{L}=\big(K,\{1,2,\cdots,N\},\{F_i\}_{i=1}^N\big)$ with $N\geq 2$. Let $\mathcal{P}_\mathcal{L}$ be its post critical set, see the general description for $\mathcal{L}$ above Proposition 2.4. Let

$V_0 = \Lambda(\mathcal{P}_{\mathcal{L}})$ be the set of boundary of K, and $V_1 = \bigcup_{i=1}^N F_i V_0$. Assume that $\mathcal{G}$ is a finite group of homeomorphisms $K \to K$, such that $g(V_0) = V_0, \forall g \in \mathcal{G}$. In addition, we require that for any $g \in \mathcal{G}$ and $1 \le i \le N$, there exists $g' \in \mathcal{G}$ and $1 \le i' \le N$ such that

$$g \circ F_i = F_{i'} \circ g'.$$

Definition 3.4. *Let $\mathcal{J}$ be an equivalent relation on V_0.*

(a). We define $\mathcal{J}^{(1)}$ to be the smallest equivalent relation on V_1 such that

$$x\mathcal{J}y \Longrightarrow F_i(x)\mathcal{J}^{(1)}F_i(y), \quad 1 \le i \le N.$$

(b). We call $\mathcal{J}$ a preserved relation if for any $x, y \in V_0$,

$$x\mathcal{J}y \Longleftrightarrow x\mathcal{J}^{(1)}y.$$

In addition, if

$$x\mathcal{J}y \Longrightarrow g(x)\mathcal{J}g(y), \quad \forall g \in \mathcal{G},$$

we call $\mathcal{J}$ a preserved $\mathcal{G}$-relation.

Remark. (a). For two equivalent relation $\mathcal{J}$, $\mathcal{J}'$ on a set V, we say $\mathcal{J}$ is smaller than $\mathcal{J}'$, denoted as $\mathcal{J} \subset \mathcal{J}'$, if each equivalent class of $\mathcal{J}$ is contained in an equivalent class of $\mathcal{J}'$.

(b). We say that $\mathcal{J}$ is *non-trivial* if $\mathcal{J}$ is neither the full relation ($x\mathcal{J}y$, $\forall x, y \in V_0$), denoted by $\mathcal{J} = 1$, nor the null relation ($x \not\mathcal{J} y$ if $x \ne y$), denoted by $\mathcal{J} = 0$.

A resistance form on a finite set V can always be written as

$$\mathcal{D}(f) = \sum_{x \ne y} j_{x,y}(f(x) - f(y))^2, \tag{3.3}$$

with nonnegative constants $j_{x,y}$. Noticing that by definition, a resistance form is non-degenerate, i.e. we always have $\mathcal{D}(f) = 0$ if and only if f is a constant function. There are also *degenerate forms* of the form (3.3) whose kernel is larger than the space of constant functions. Clearly, $\mathcal{D}$ is non-degenerate if and only if the matrix $(j_{x,y})_{x,y \in V}$ is irreducible.

Definition 3.5. *(a). Let $\mathcal{J}$ be a preserved relation on V_0, and $\mathcal{D}$ be a form of the form (3.3).*

(a1). We say $\mathcal{D} \in \mathcal{M}_{\mathcal{J}}$ if

$$\mathcal{D}(f) = 0 \Longleftrightarrow f \text{ is constant on each equivalent class of } \mathcal{J}.$$

(a2). We define $T_{\mathcal{J}} : \mathcal{M}_{\mathcal{J}} \to \mathcal{M}_{\mathcal{J}}$ as follows,

$$T_{\mathcal{J}}\mathcal{D}(f) = \inf\{\mathcal{D}^{(1)}(\tilde{f}) : \tilde{f} = f \text{ on } V_0, \tilde{f} \in l(V_1)\}, \quad \forall f \in l(V_0),$$

where $\mathcal{D}^{(1)}(\tilde{f}) = \sum_{i=1}^N \mathcal{D}(\tilde{f} \circ F_i)$.

(b). Let $\mathcal{M}_{V_0/\mathcal{J}}$ be the space of resistance forms on $V_0/\mathcal{J}$. We identify $l(V_0/\mathcal{J})$ with the subspace of $l(V_0)$ where each f admits constant values on each equivalent class of $\mathcal{J}$.

(b1). For each resistance form $\mathcal{D}$ on V_0, we can naturally define a form $\mathcal{D}_{V_0/\mathcal{J}}$ in $\mathcal{M}_{V_0/\mathcal{J}}$ by

$$\mathcal{D}_{V_0/\mathcal{J}}(f) = \mathcal{D}(f), \quad \forall f \in l(V_0/\mathcal{J}).$$

In addition, any form in $\mathcal{M}_{V_0/\mathcal{J}}$ can be constructed in this way.

(b2). We define $T_{V_0/\mathcal{J}} : \mathcal{M}_{V_0/\mathcal{J}} \to \mathcal{M}_{V_0/\mathcal{J}}$ as follows,

$$T_{V_0/\mathcal{J}}\mathcal{D}_{V_0/\mathcal{J}}(f) = \inf\{\mathcal{D}^{(1)}(\tilde{f}) : \tilde{f} = f \text{ on } V_0, \tilde{f} \in l(V_1/\mathcal{J}^{(1)})\}, \quad \forall f \in l(V_0/\mathcal{J}),$$

where $\mathcal{D}^{(1)}(\tilde{f}) = \sum_{i=1}^N \mathcal{D}(\tilde{f} \circ F_i)$.

Definition 3.6. *Let $\mathcal{J}$ be a preserved relation on V_0.*

(a). We define

$$\underline{\rho}_{\mathcal{J}}(\mathcal{D}) = \inf_{f\in l(V_0)\setminus l(V_0/\mathcal{J})} \frac{T_{\mathcal{J}}\mathcal{D}(f)}{\mathcal{D}(f)}, \quad \overline{\rho}_{\mathcal{J}}(\mathcal{D}) = \sup_{f\in l(V_0)\setminus l(V_0/\mathcal{J})} \frac{T_{\mathcal{J}}\mathcal{D}(f)}{\mathcal{D}(f)}, \text{ for } \mathcal{D}\in\mathcal{M}_{\mathcal{J}},$$

and

$$\underline{\rho}_{V_0/\mathcal{J}}(\mathcal{D}) = \inf_{f\in l(V_0/\mathcal{J})} \frac{T_{V_0/\mathcal{J}}\mathcal{D}(f)}{\mathcal{D}(f)}, \quad \overline{\rho}_{V_0/\mathcal{J}}(\mathcal{D}) = \sup_{f\in l(V_0/\mathcal{J})} \frac{T_{V_0/\mathcal{J}}\mathcal{D}(f)}{\mathcal{D}(f)}, \text{ for } \mathcal{D}\in\mathcal{M}_{V_0/\mathcal{J}}.$$

(b). We define

$$\underline{\rho}_{\mathcal{J}} = \sup_{\mathcal{D}\in\mathcal{M}_{\mathcal{J}}} \underline{\rho}_{\mathcal{J}}(\mathcal{D}), \quad \overline{\rho}_{\mathcal{J}} = \inf_{\mathcal{D}\in\mathcal{M}_{\mathcal{J}}} \overline{\rho}_{\mathcal{J}}(\mathcal{D}),$$

$$\underline{\rho}_{V_0/\mathcal{J}} = \sup_{\mathcal{D}\in\mathcal{M}_{V_0/\mathcal{J}}} \underline{\rho}_{V_0/\mathcal{J}}(\mathcal{D}), \quad \overline{\rho}_{V_0/\mathcal{J}} = \inf_{\mathcal{D}\in\mathcal{M}_{V_0/\mathcal{J}}} \overline{\rho}_{V_0/\mathcal{J}}(\mathcal{D}).$$

(c). If $\mathcal{J}$ is in addition a preserved $\mathcal{G}$-relation, we define $\underline{\rho}^{\mathcal{G}}_{\mathcal{J}} = \sup_{\mathcal{D}} \underline{\rho}_{\mathcal{J}}(\mathcal{D})$ where the supremum is taken over $\mathcal{G}$-symmetric forms in $\mathcal{M}_{\mathcal{J}}$. In particular, when $\mathcal{G}$ is taken to be the trivial group, i.e. $\mathcal{G} = \{id\}$, then $\underline{\rho}^{\mathcal{G}}_{\mathcal{J}} = \underline{\rho}_{\mathcal{J}}$.

$\overline{\rho}^{\mathcal{G}}_{\mathcal{J}}$, $\underline{\rho}^{\mathcal{G}}_{V_0/\mathcal{J}}$ and $\overline{\rho}^{\mathcal{G}}_{V_0/\mathcal{J}}$ are defined in a same way.

We now quote the theorem of Sabot which will serve as our main instrument for proving the existence and uniqueness.

Theorem 3.7 ([35])**.** *(a). If there exist two non-trivial preserved $\mathcal{G}$-relations $\mathcal{J}$ and $\mathcal{J}'$ on V_0, such that $\underline{\rho}^{\mathcal{G}}_{V_0/\mathcal{J}} < \underline{\rho}^{\mathcal{G}}_{\mathcal{J}'}$, then (3.2) does not have a $\mathcal{G}$-symmetric solution.*

(b). If for all non-trivial preserved $\mathcal{G}$-relation $\mathcal{J}$, it holds that $\overline{\rho}^{\mathcal{G}}_{\mathcal{J}} < \underline{\rho}^{\mathcal{G}}_{V_0/\mathcal{J}}$, then (3.2) has at most one $\mathcal{G}$-symmetric solution (up to a multiplicative constant). If moreover, there do not exist two strictly ordered non-trivial $\mathcal{G}$-relations (i.e. $\mathcal{J}\subset\mathcal{J}'$ and $\mathcal{J}\neq\mathcal{J}'$), then we have exactly one $\mathcal{G}$-symmetric solution to (3.2).

Remark. For the uniqueness part, the inequality in Theorem 3.7 can be loosed to

$$\overline{\rho}^{\mathcal{G}}_{\mathcal{J},k} < \underline{\rho}^{\mathcal{G}}_{V_0/\mathcal{J},k} \quad \text{for some } k\in\mathbb{N}.$$

Here $\overline{\rho}^{\mathcal{G}}_{\mathcal{J},k} = \inf_{\mathcal{D}\in\mathcal{M}_{\mathcal{J}}} \sup_{f\in l(V_0)\setminus l(V_0/\mathcal{J})} \frac{T^k_{\mathcal{J}}\mathcal{D}(f)}{\mathcal{D}(f)}$ and $\underline{\rho}^{\mathcal{G}}_{V_0/\mathcal{J},k} = \sup_{\mathcal{D}\in\mathcal{M}_{V_0/\mathcal{J}}} \inf_{f\in l(V_0/\mathcal{J})} \frac{T^k_{V_0/\mathcal{J}}\mathcal{D}(f)}{\mathcal{D}(f)}$, where the supremum and infimum are taken over $\mathcal{G}$-symmetric forms. Indeed, readers can revise Lemma 5.7 in Sabot's paper [35] with this new assumption. The rest of the proof of the uniqueness in Section 5.4 of [35] then follows in the same way.

3.2. **Proof of existence.** We return to the study of the Julia sets K_λ associated with the MS maps of the form $R_\lambda(z) = z^n + \frac{\lambda}{z^m}$ with $n\geq 2, m\geq 1$ and $\lambda\in\mathbb{C}\setminus\{0\}$. In this subsection, we will prove the existence of a balanced resistance form on K_λ. By the discussion in Section 3.1, it is enough to study the equation (3.2) by applying Sabot's theorem. However, since the fractal depends in a complicated manner on R_λ, we must be careful in verifying the conditions in Theorem 3.7.

Throughout this subsection, we will let $\mathcal{G}$ be the canonical rotation group on K_λ, that is $\mathcal{G} = \{g : g(x) = e^{\frac{2l\pi i}{m+n}}x, 0\leq l\leq m+n-1\}$. We first note several properties of preserved relations on V_0.

Definition 3.8. *Let $G=(V,E)$ and $G'=(V',E')$ be two finite graphs, and $f:V\to V'$.*
(a). We define $f(G)=(f(V),f(E))$ to be the graph with vertices $f(V)$ and

$$f(E)=\big\{\{f(x),f(y)\}:\{x,y\}\in E, f(x)\neq f(y)\big\}.$$

(b). We define $G\cup G'$ to be the graph with vertices $V\cup V'$ and edges $E\cup E'$.

Definition 3.9. *Let $\mathcal{J}$ be a preserved relation on V_0.*
(a). Define $G_{\mathcal{J}}=(V_0,E_{\mathcal{J}})$ be a graph with vertices V_0, and $\{x,y\}\in E_{\mathcal{J}}$ if and only if $x\mathcal{J}y$.
(b). For $k\geq 0$, we define $V_k=\bigcup_{|w|=k}F_wV_0$ and the graph $G^{(k)}_{\mathcal{J}}(=(V_k,E^{(k)}_{\mathcal{J}}))=\bigcup_{|w|=k}F_kG_{\mathcal{J}}$. Here we use w to represent a finite word and $|w|$ to denote its length, so we sum over $w\in\{1,2,\cdots,m+n\}^k$. The notation F_w is short for $F_{w_1}\circ F_{w_2}\circ\cdots\circ F_{w_k}$.
(c). We define an equivalent relation $\mathcal{J}^{(k)}$ on V_k by

$$x\mathcal{J}^{(k)}y\Longleftrightarrow x \text{ and } y \text{ belong to the same connected component of } G^{(k)}_{\mathcal{J}}.$$

For a sequence $x=x_0,x_1,x_2,\cdots,x_L=y$ such that $\{x_{i-1},x_i\}\in E^{(k)}_{\mathcal{J}}$, we call it a $G^{(k)}_{\mathcal{J}}$-path connecting x and y.
We can also consider a preserved relation $\mathcal{J}$ on $\tilde{V}_0$, and define $\tilde{G}_{\mathcal{J}}$, $\tilde{G}^{(k)}_{\mathcal{J}}$ and $\mathcal{J}^{(k)}$ in the corresponding way.

The following lemma, which shows the relation between $\mathcal{J}^{(k)},k\geq 0$, will play a fundamental role throughout this section. In particular, (b), (c) are special properties of the Julia sets, and do not hold for general p.c.f. fractals.

Lemma 3.10. *Let $\mathcal{J}$ be a preserved relation on V_0 (or $\tilde{V}_0$). Then the definitions of $\mathcal{J}^{(1)}$ in Definition 3.4 (a) and Definition 3.9 (c) coincide. More generally, for any $0\leq k<l$, we have $\mathcal{J}^{(l)}$ is the smallest equivalent relation such that*

$$x\mathcal{J}^{(k)}y\Longrightarrow F_w(x)\mathcal{J}^{(l)}F_w(y),\quad \forall|w|=l-k \text{ and } x,y\in V_k \text{ (or } \tilde{V}_k).$$

In addition, we have:
(a). Let $0\leq k<l$, then for any $x,y\in V_k$ (or $\tilde{V}_k$),

$$x\mathcal{J}^{(k)}y\Longleftrightarrow x\mathcal{J}^{(l)}y.$$

(b). Let $k\geq 1$, then for any $x,y\in V_k$ (or $\tilde{V}_k$),

$$x\mathcal{J}^{(k)}y\Longrightarrow R_\lambda(x)\mathcal{J}^{(k-1)}R_\lambda(y).$$

(c). Let $k\geq 0$, then for any $x,y\in V_k$ (or $\tilde{V}_k$),

$$x\mathcal{J}^{(k)}y\Longrightarrow R_\lambda(x)\mathcal{J}^{(k)}R_\lambda(y).$$

Proof. Assume $x\mathcal{J}^{(k)}y$, then there exists a $G^{(k)}_{\mathcal{J}}$-path $x=x_0,x_1,x_2,\cdots,x_L=y$ connecting x and y. Clearly, for any $|w|=l-k$, we have that $F_wx_0,F_wx_1,\cdots,F_wx_L$ is a $G^{(l)}_{\mathcal{J}}$-path. This shows

$$x\mathcal{J}^{(k)}y\Longrightarrow F_w(x)\mathcal{J}^{(l)}F_w(y),\quad \forall|w|=l-k.$$

Noticing that $G^{(l)}_{\mathcal{J}}=\bigcup_{|w|=l-k}F_w(G^{(k)}_{\mathcal{J}})$, we have $E^{(l)}_{\mathcal{J}}\subset\big\{\{F_w(x),F_w(y)\}:x\mathcal{J}^{(k)}y,|w|=l-k\big\}$. Since $\mathcal{J}^{(l)}$ is generated by the edge set $E^{(l)}_{\mathcal{J}}$, we claim that $\mathcal{J}^{(l)}$ is the smallest relation such that the above implication holds.

(a). We view an equivalent relation $\mathcal{J}^{(k)}$ as a subset of $V_k \times V_k$, i.e. for $(x,y) \in V_k \times V_k$, $(x,y) \in \mathcal{J}^{(k)}$ if and only if $x\mathcal{J}^{(k)}y$. Note that $\{V_k \times V_k\}_{k\geq 0}$ is an increasing sequence of sets. Then, we have $\mathcal{J} \subset \mathcal{J}^{(1)}$ as $\mathcal{J}$ is preserved. Noticing that $\mathcal{J}^{(1)}$ is generated with $\mathcal{J}$, and $\mathcal{J}^{(2)}$ is generated with $\mathcal{J}^{(1)}$ in a same manner, we have $\mathcal{J}^{(1)} \subset \mathcal{J}^{(2)}$. Continuing the argument, we get

$$\mathcal{J} \subset \mathcal{J}^{(1)} \subset \mathcal{J}^{(2)} \subset \mathcal{J}^{(3)} \subset \cdots .$$

This shows that for any $x, y \in V_k$, $x\mathcal{J}^{(k)}y \Longrightarrow x\mathcal{J}^{(l)}y$.

For the other direction, we assume that $x, y \in V_k$ and $x\mathcal{J}^{(l)}y$. Then there exists a $G_{\mathcal{J}}^{(l)}$-path $x = x_0, x_1, x_2, \cdots, x_L = y$ connecting x and y. We choose a subsequence $x = x_{i_0}, x_{i_1}, \cdots, x_{i_M} = y$ such that $0 = i_0 < i_1 < i_2 < \cdots < i_M = L$, and

$$\{i_k\}_{k=1}^{M-1} = \{1 \leq j \leq L-1 : x_j \in V_{l-1}\}.$$

Now, we look at x_{i_0} and x_{i_1}. Clearly, $x_{i_0}, x_{i_0+1}, x_{i_0+2}, \cdots, x_{i_1}$ is a $G_{\mathcal{J}}^{(l)}$-path, which is contained in a same $(l-1)$-cell $F_w K_\lambda$. So $F_w^{-1}x_{i_0}, F_w^{-1}x_{i_0+1}, \cdots, F_w^{-1}x_{i_1}$ is a $G_{\mathcal{J}}^{(1)}$-path, thus we have $(F_w^{-1}x_{i_0})\mathcal{J}^{(1)}(F_w^{-1}x_{i_1})$, and so $(F_w^{-1}x_{i_0})\mathcal{J}(F_w^{-1}x_{i_1})$. This implies $\{x_{i_0}, x_{i_1}\} \in E_{\mathcal{J}}^{(l-1)}$. By the same argument, we see that $\{x_{i_1}, x_{i_2}\} \in E_{\mathcal{J}}^{(l-1)}$, $\{x_{i_2}, x_{i_3}\} \in E_{\mathcal{J}}^{(l-1)}$, $\cdots$. So, $x = x_{i_0}, x_{i_1}, x_{i_2}, \cdots, x_{i_M} = y$ is a $G_{\mathcal{J}}^{(l-1)}$-path, and $x\mathcal{J}^{(l-1)}y$. By repeating the above arguments, we have

$$x\mathcal{J}^{(l)}y \Longrightarrow x\mathcal{J}^{(l-1)}y \Longrightarrow \cdots \Longrightarrow x\mathcal{J}^{(k)}y.$$

(b). Let $x\mathcal{J}^{(k)}y$, then there is a $G_{\mathcal{J}}^{(k)}$-path $x, x_1, x_2, \cdots, x_L = y$. Then we have $R_\lambda(x)$, $R_\lambda(x_1), \cdots, R_\lambda(y)$ is a $G_{\mathcal{J}}^{(k-1)}$-path, noticing that $G_{\mathcal{J}}^{(k)} = \bigcup_{i=1}^{m+n} F_i G_{\mathcal{J}}^{(k-1)}$. Thus, this implies that $R_\lambda(x)\mathcal{J}^{(k-1)}R_\lambda(y)$.

(c) This assertion is an easy synthesis of (a) and (b). We have for $x, y \in V_k$, $k \geq 0$,

$$x\mathcal{J}^{(k)}y \Longleftrightarrow x\mathcal{J}^{(k+1)}y \Longrightarrow R_\lambda(x)\mathcal{J}^{(k)}R_\lambda(y).$$

Finally, we point out that the proof for the $\tilde{V}_0$ setting is the same. □

As an important corollary to the Lemma 3.10, we have the following lemma concerning the critical set C.

Lemma 3.11. *(a). Let $\mathcal{J}$ be a preserved relation on V_0 and assume $x\mathcal{J}^{(1)}y$ for any $x, y \in C$, then $\mathcal{J} = 1$.*

(b). Let $\mathcal{J}$ be a preserved $\mathcal{G}$-relation on $\tilde{V}_0$ and assume $x\mathcal{J}^{(1)}y$ for any $x, y \in C$, then $\mathcal{J} = 1$.

(c). Let $\mathcal{J}$ be a non-trivial preserved $\mathcal{G}$-relation on $\tilde{V}_0$, we have $x\not\mathcal{J}^{(1)}y$ for any distinct $x, y \in C$.

Proof. (a). By Lemma 3.10 (c), we have $x\mathcal{J}^{(1)}y$ for any distinct $x, y \in R_\lambda(C)$. Fix a distinct pair of $x, y \in R_\lambda(C)$. Note that due to Lemma 2.1 (a) and the fact $e^{\frac{2\pi i}{m+n}}C = C$, we have $y = e^{\frac{2k\pi i}{m+n}}x$ for some $k \in \mathbb{Z}, e^{\frac{2k\pi i}{m+n}} \neq 1$. Then, by (2.3), $x \in F_i K_\lambda, y \in F_j K_\lambda$ for some $1 \leq i \leq n+m$ and $j \equiv i + k \mod n+m$. Note that x, y do not live in other 1-cells, since by Lemma 2.1 (b), $x, y \notin C = \bigcup_{i \neq j} F_i K_\lambda \cap F_j K_\lambda$. So we conclude that x, y can not belong to a same 1-cell of K_λ. Let $x, x_1, x_2, \cdots, y$ be a $G_{\mathcal{J}}^{(1)}$-path connecting x and y, then the path must exit the 1-cell $F_i K_\lambda$ containing x at some point $x_l \in C \cap F_i V_0$ by Corollary 2.3 and equation (2.1) (also see Figure 3 for an illustration). As a consequence, we have $x\mathcal{J}^{(1)}x_l$ with $x \in R_\lambda(C), x_l \in C$. This implies that

$$x\mathcal{J}^{(1)}y, \quad \forall x, y \in R_\lambda(C) \cup C,$$

noticing that we already have $x\mathcal{J}^{(1)}y$ for any $x, y \in C$ and $x\mathcal{J}^{(1)}y$ for any $x, y \in R_\lambda(C)$. In addition, by using Lemma 3.10 (c) again, for any $k \geq 0$, we still have

$$x\mathcal{J}^{(1)}y, \quad \forall x, y \in R_\lambda^k\big(R_\lambda(C) \cup C\big).$$

Taking the union over k, we see that

$$x\mathcal{J}^{(1)}y, \quad \forall x, y \in \bigcup_{k=0}^{\infty} R_\lambda^k(C).$$

This implies

$$x\mathcal{J}y, \quad \forall x, y \in V_0 = \bigcup_{k=1}^{\infty} R_\lambda^k(C),$$

so $\mathcal{J} = 1$.

(b). By a same proof of (a), one can see that $x\mathcal{J}^{(1)}y$ for any $x, y \in V_0 \cup C$. Since $\mathcal{J}^{(1)}$ is a $\mathcal{G}$-relation, we have $x\mathcal{J}^{(1)}y$ for any $x, y \in \tilde{V}_0 \cup C$. Thus $\mathcal{J} = 1$ as desired.

(c). For the sake of contradiction we assume there exists $x \neq y$ in C such that $x\mathcal{J}^{(1)}y$, and consider two cases.

Case 1: $y = e^{\frac{2\pi i}{m+n}}x$. In this case, by rotation symmetry ($\mathcal{G}$-symmetry), we have

$$x\mathcal{J}^{(1)}e^{\frac{2\pi i}{m+n}}x\mathcal{J}^{(1)}e^{\frac{4\pi i}{m+n}}x\mathcal{J}^{(1)}\cdots,$$

so C is in a same class of $\mathcal{J}^{(1)}$. This implies $\mathcal{J} = 1$ by (b), and gives a contradiction.

Case 2: $y = e^{\frac{2k\pi i}{m+n}}x$ for some $2 \leq k \leq m+n-2$. In this case, any $G_{\mathcal{J}}^{(1)}$-path $x, x_1, x_2, \cdots, y$ contains a vertice in $\{e^{\pm\frac{2\pi i}{m+n}}x\} \subset C$, since removing the vertices $\{e^{\pm\frac{2\pi i}{m+n}}x\}$ will disconnect x and $C \setminus \{x, e^{\pm\frac{2\pi i}{m+n}}x\}$. This implies that we can find $y' = e^{\frac{2\pi i}{m+n}}x$ or $e^{-\frac{2\pi i}{m+n}}x$ in C such that $x\mathcal{J}^{(1)}y'$, which reduces the problem to Case 1. □

Lemma 3.11 shows that for a non-trivial preserved $\mathcal{G}$-relation $\mathcal{J}$, the extension $\mathcal{J}^{(1)}$ is quite loose. In particular, there are few choices of paths for x, y in different 1-cells.

Lemma 3.12. *Let $\mathcal{J}$ be a non-trivial preserved $\mathcal{G}$-relation on $\tilde{V}_0$, and assume $x, y \in \tilde{V}_1 \setminus C$ with $x\mathcal{J}^{(1)}y$. Then there exists $1 \leq i \leq m+n$ such that $\{x, y\} \subset F_i\tilde{V}_0 \cup F_{i+1}\tilde{V}_0$ (cyclic notation $m+n+1 = 1$). In addition, if x, y belong to different 1-cells, say $x \in F_i\tilde{V}_0$, $y \in F_{i+1}\tilde{V}_0$. Then we have $x\mathcal{J}^{(1)}c_i\mathcal{J}^{(1)}y$ (Recall (2.2)).*

Proof. Let $x = x_0, x_1, \cdots, x_N = y$ be a $G_{\mathcal{J}}^{(1)}$-path connecting x and y. Assume $x \in F_i\tilde{V}_0$ and $y \in F_j\tilde{V}_0$. If $i \neq j$, then the path should leave $F_i\tilde{V}_0$ at c_{i-1} or c_i, and enter $F_j\tilde{V}_0$ at c_{j-1} or c_j. However, according to Lemma 3.11 (c), there is at most one critical point contained in the path (not counting multiplicity). This is only possible when $|j - i| = 1$.

If $x \in F_i\tilde{V}_0$ and $y \in F_{i+1}\tilde{V}_0$, then any path from x to y includes at least one $c \in C$, and the only possible choice is c_i as discussed above. □

In the rest of this section, we will prove the existence of a $\mathcal{G}$-symmetric solution to (3.2). We consider the cases $m \geq 2$ and $m = 1$ separately since they exhibit quite different properties.

- **The $m \geq 2$ case.**

Throughout this part, we assume that R_λ takes the form $R_\lambda(z) = z^n + \frac{\lambda}{z^m}$ with $n \geq 2, m \geq 2$. In this case, we will show that *there are only trivial preserved $\mathcal{G}$-relations.* To use the dynamics more efficiently, we introduce a natural distance on β_λ. Recall that β_λ is the boundary of the immediate attracting basin of ∞, and $R_\lambda : \beta_\lambda \to \beta_\lambda$ is conjugate to $\Phi_n : \mathbb{T} \to \mathbb{T}$ via a homeomorphism $\psi_\lambda : \beta_\lambda \to \mathbb{T}$ as illustrated in Lemma 2.5.

Definition 3.13. *Let $x, y \in \beta_\lambda$, we define*

$$d_{\beta_\lambda}(x, y) = d_{\mathbb{T}}(\psi_\lambda(x), \psi_\lambda(y)),$$

where $d_{\mathbb{T}}$ is the standard distance on the unit circle $\mathbb{T}$ (for $a, b \in [0, 1)$, we have $d_{\mathbb{T}}([a], [b]) = \min\{|a - b|, 1 - |a - b|\}$).

By the conjugacy of the dynamics of R_λ and Φ_n on β_λ, it is easy to see the following result.

Lemma 3.14. *Let $x, y \in \beta_\lambda$, we have $d_{\beta_\lambda}(R_\lambda(x), R_\lambda(y)) = \min\{n d_{\beta_\lambda}(x, y), 1 - n d_{\beta_\lambda}(x, y)\}$ if $d_{\beta_\lambda}(x, y) < \frac{1}{n}$.*

The distance on β_λ allows us to conveniently show when two vertices in $\tilde{V}_0$ are in distinct $\mathcal{J}^{(1)}$ classes, noting that $\tilde{V}_0 \subset \beta_\lambda$.

Lemma 3.15. *Let $\mathcal{J}$ be a non-trivial preserved $\mathcal{G}$-relation on $\tilde{V}_0$ and $x, y \in \tilde{V}_0$. We have*

$$d_{\beta_\lambda}(x, y) \geq \frac{2}{n(m + n)} \Longrightarrow x \not\mathcal{J} y.$$

Proof. Let $x, y \in \tilde{V}_0$ with $d_{\beta_\lambda}(x, y) \geq \frac{2}{n(m+n)}$. By Lemma 3.10 (a) and Lemma 3.12, without loss of generality, we may assume that for some i, $x \in F_i\tilde{V}_0$ and $y \in F_i\tilde{V}_0 \cup F_{i+1}\tilde{V}_0$. Otherwise, we already have $x \not\mathcal{J} y$. We consider two cases.

Case 1: $y \in F_i\tilde{V}_0$. In this case, x and y belong to a same 1-cell, so $d_{\beta_\lambda}(x, y) < \frac{1}{m+n}$. Consequently, by Lemma 3.14 and the condition $d_{\beta_\lambda}(x, y) \geq \frac{2}{n(m+n)}$, we have

$$d_{\beta_\lambda}(R_\lambda(x), R_\lambda(y)) \geq \min\{n \cdot \frac{2}{n(m + n)}, 1 - n \cdot \frac{1}{m + n}\} = \frac{2}{m + n},$$

noting that in the last equality we use the condition $m \geq 2$. This means that $R_\lambda(x)$ and $R_\lambda(y)$ do not belong to neighbouring 1-cells, so by Lemma 3.12, $R_\lambda(x) \not\mathcal{J}^{(1)} R_\lambda(y)$. Finally, we apply Lemma 3.10 to see $x \not\mathcal{J} y$.

Case 2: $y \in F_{i+1}\tilde{V}_0$. We prove $x \not\mathcal{J} y$ by contradiction. Assume that $x\mathcal{J}y$, then by Lemma 3.12, we have

$$x \mathcal{J}^{(1)} c_i \mathcal{J}^{(1)} y.$$

See Figure 6 for an illustration. We label the two vertices in $R_\lambda^{-1}(C) \cap \beta_\lambda$ surrounding c_i by $\tilde{c}_1, \tilde{c}_2$, ordered so that $\tilde{c}_1 \in F_i\tilde{V}_0$ and $\tilde{c}_2 \in F_{i+1}\tilde{V}_0$.

Case 2.1: c_i, x belong to the same 2-cell. In this case, we claim that c_i and y do not belong to neighbouring 2-cells. If they do, then $d_{\beta_\lambda}(x, c_i) < d_{\beta_\lambda}(c_i, \tilde{c}_1)$ and by Lemma 2.7, $d_{\beta_\lambda}(y, c_i) < d_{\beta_\lambda}(c_i, \tilde{c}_2) + \frac{1}{n(m+n)}$, which implies that $d_{\beta_\lambda}(x, y) \leq d_{\beta_\lambda}(x, c_i) + d_{\beta_\lambda}(y, c_i) < \frac{2}{n(m+n)}$, noticing

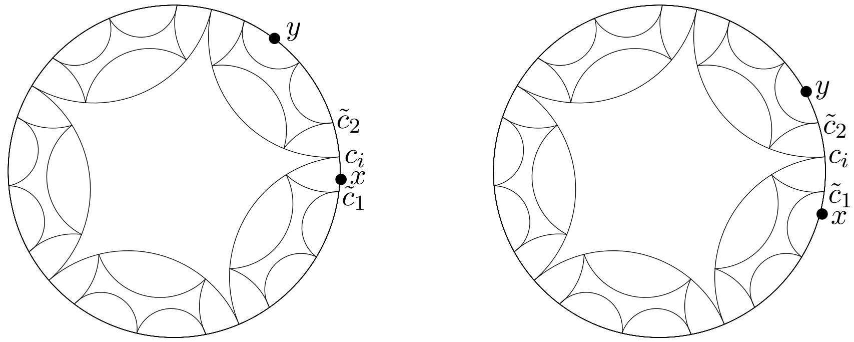


FIGURE 6. An illustration of the vertices $x, \tilde{c}_1, c_i, \tilde{c}_2, y$.

that $d_{\beta_\lambda}(c_i, \tilde{c}_1) + d_{\beta_\lambda}(c_i, \tilde{c}_2) = \frac{1}{n(m+n)}$, a contradiction. It is then easy to see that $R_\lambda(c_i)$ and $R_\lambda(y)$ do not belong to neighbouring 1-cells, and thus by Lemma 3.12, $R_\lambda(c_i) \not\mathcal{J}^{(1)} R_\lambda(y)$. Then using Lemma 3.10 (c), we have $c_i \not\mathcal{J}^{(1)} y$, violating the initial assumption of Case 2.

Case 2.2: c_i, x belong to distinct 2-cells. We can additionally assume that c_i, y also belong to distinct 2-cells, otherwise we are essentially back to Case 2.1. Clearly by Lemma 3.10 (c), we have

$$x\mathcal{J}^{(1)}c_i \Longrightarrow R_\lambda(x)\mathcal{J}^{(1)}R_\lambda(c_i).$$

Then by Lemma 3.12 and Lemma 3.14, $R_\lambda(x)$ and $R_\lambda(c_i)$ must belong to two neighbouring 1-cells separately which intersection at $R_\lambda(\tilde{c}_1)$. This gives that $R_\lambda(\tilde{c}_1)\mathcal{J}^{(1)}R_\lambda(c_i)$. By the symmetric argument we have $R_\lambda(\tilde{c}_2)\mathcal{J}^{(1)}R_\lambda(c_i)$. This implies $R_\lambda(\tilde{c}_1)\mathcal{J}^{(1)}R_\lambda(\tilde{c}_2)$, contradicting Lemma 3.11 (c), noticing that $R_\lambda(\tilde{c}_1), R_\lambda(\tilde{c}_2) \in C$. □

Remark. The above proof indirectly uses $\mathcal{J}^{(2)}$. More specifically, the proof of $x\mathcal{J}^{(1)}y \Longrightarrow R_\lambda(x)\mathcal{J}^{(1)}R_\lambda(y)$ in Lemma 3.10 (c) essentially involves going to the second level.

By applying Lemma 3.14 and 3.15, we can finally prove the non-existence of non-trivial preserved $\mathcal{G}$-relation when $m \geq 2$.

Proposition 3.16. *Let $R_\lambda(z) = z^n + \frac{\lambda}{z^m}$ with $m, n \geq 2$ be a MS map. There does not exist a non-trivial preserved $\mathcal{G}$-preserved relation on $\tilde{V}_0$. In particular, there exists exactly one $\mathcal{G}$-symmetric solution to (3.2).*

Proof. We prove by contradiction. Let $\mathcal{J}$ be a non-trivial preserved $\mathcal{G}$-relation. Take $x \neq y$ from $\tilde{V}_0$. We claim that there exists $k \geq 0$ such that

$$d_{\beta_\lambda}(R_\lambda^k(x), R_\lambda^k(y)) \geq \frac{2}{n(m+n)}. \tag{3.4}$$

In fact, if not, then for all $k \geq 0$, we have $d_{\beta_\lambda}(R_\lambda^k(x), R_\lambda^k(y)) < \frac{2}{n(m+n)}$, which implies

$$1 - nd_{\beta_\lambda}(R_\lambda^{k-1}(x), R_\lambda^{k-1}(y)) > 1 - n\frac{2}{n(m+n)} \geq \frac{2}{n(m+n)} > d_{\beta_\lambda}(R_\lambda^k(x), R_\lambda^k(y)), \quad \forall k \geq 1.$$

Thus, by Lemma 3.14,

$$d_{\beta_\lambda}(R_\lambda^k(x), R_\lambda^k(y)) = nd_{\beta_\lambda}(R_\lambda^{k-1}(x), R_\lambda^{k-1}(y)) = \cdots = n^k d_{\beta_\lambda}(x, y).$$

Letting $k \to \infty$, we get $d_{\beta_\lambda}(x,y) = 0$, a contradiction.

Now, to prove existence for the $m \geq 2$ case, choose $k \geq 0$ such that (3.4) holds. Then, by Lemma 3.15, we have $R_\lambda^k(x) \not\mathcal{J} R_\lambda^k(y)$. Thus, applying Lemma 3.10 (c), we have $x \not\mathcal{J} y$. Noticing that x, y are arbitrarily chosen, we have $\mathcal{J} = 0$. A contradiction. Thus there does not exist a non-trivial preserved $\mathcal{G}$-relation on $\tilde{V}_0$.

Finally, by applying Sabot's theorem, Theorem 3.7 (b), there exists exactly one $\mathcal{G}$-symmetric solution to (3.2). □

- **The $m = 1$ case.**

In this case, there may exist a non-trivial preserved $\mathcal{G}$-relation on $\tilde{V}_0$ $(= V_0)$.

Example 3.17. *We consider the first Julia set presented in Example 2.6. We define an equivalent relation $\mathcal{J}$ by taking each pair of 'opposite' vertices to be an equivalent class. More precisely, there are three class I_1, I_2, I_3 in $\mathcal{J}$, with*

$$I_1 = \psi_\lambda^{-1}\{[0], [\frac{1}{2}]\}, \quad I_2 = \psi_\lambda^{-1}\{[\frac{1}{6}], [\frac{2}{3}]\}, \quad I_3 = \psi_\lambda^{-1}\{[\frac{1}{3}], [\frac{5}{6}]\}.$$

See Figure 7 for an illustration of $\mathcal{J}$ and $\mathcal{J}^{(1)}$. On can see that $\mathcal{J}$ is a perserved $\mathcal{G}$-relation.

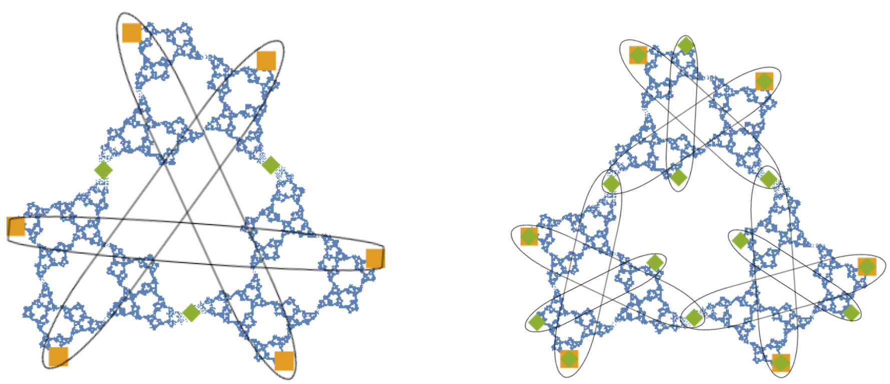


FIGURE 7. An illustration of a non-trivial preserved $\mathcal{G}$-relation $\mathcal{J}$

.

In the subsequent lemmas, we provide a rough picture of all possible preserved $\mathcal{G}$-relations. The proof of the $m \geq 2$ case does not work here, but a similar argument still provides us some insight. Also, as $m+n$ and n are coprime (now $m = 1$), we can easily see that $\tilde{V}_0 = V_0$. We substitute $m = 1$ in the following discussions.

Lemma 3.18. *Let $\mathcal{J}$ be a preserved relation on V_0 such that*

$$x \not\mathcal{J}^{(1)} y, \quad \forall x, y \in C \cup R_\lambda(C).$$

Then we have $\mathcal{J} = 0$.

Proof. The proof of the lemma is similar to that of Proposition 3.16. In fact, we will show the following claim analogous to Lemma 3.15.

Claim: $d_{\beta_\lambda}(x,y) \geq \frac{1}{n+1} \Longrightarrow x \not\mathcal{J} y$ for $x, y \in V_0$.

Proof of the claim. Let $x, y \in V_0$ with $d_{\beta_\lambda}(x,y) \geq \frac{1}{n+1}$. We will prove $x\not\mathcal{J}y$ by contradiction. Suppose $x\mathcal{J}y$. Clearly, x, y belong to different 1-cells, and by Lemma 3.10 (a) and Lemma 3.12 there is $1 \leq i \leq n+1$ such that $x \in F_iV_0$, $y \in F_{i+1}V_0$ and $x\mathcal{J}^{(1)}c_i\mathcal{J}^{(1)}y$. In addition, one can see that $d_{\beta_\lambda}(x,c_i) + d_{\beta_\lambda}(c_i,y) = d_{\beta_\lambda}(x,y) \geq \frac{1}{n+1}$. Without loss of generality, we assume $d_{\beta_\lambda}(x,c_i) \geq \frac{1}{2(n+1)} \geq \frac{1}{n(n+1)}$, which means that x, c_i are not in the same 2-cell. Noticing that $R_\lambda(x)\mathcal{J}^{(1)}R_\lambda(c_i)$ by Lemma 3.10 (c). By Lemma 3.12, there is a $G^{(1)}_{\mathcal{J}}$-path connecting $R_\lambda(x), R_\lambda(c_i)$, which pass through some $c \in C$. Thus, $c\mathcal{J}^{(1)}R_\lambda(c_i)$, which is a contradiction to the assumption of the lemma. Hence it follows that $x\not\mathcal{J}y$.

The lemma follows from the above claim and the argument underlying Proposition 3.16, with $\frac{2}{n(m+n)}$ there replaced by $\frac{1}{n+1}$. □

By Lemma 3.11 and Lemma 3.18, we have that the restriction of a non-trivial preserved relation $\mathcal{J}^{(1)}$ to $C \cup R_\lambda(C)$ is also non-trivial. This shows that there are not too many non-trivial preserved $\mathcal{G}$-relations. To quantify this, we define two possible candidates.

Definition 3.19. *Define κ as the unique permutation on $\{1, 2, \cdots, n+1\}$ such that*

$$R_\lambda(c_{\kappa(i)}) \in F_i(K_\lambda).$$

(a). Define $\breve{G}_+ = (C \cup R_\lambda(C), \breve{E}_+)$, with the edge set

$$\breve{E}_+ = \{\{c_i, R_\lambda(c_{\kappa(i)})\} : 1 \leq i \leq n+1\}.$$

Define $G_+ = \bigcup_{k=1}^\infty R^k_\lambda(\breve{G}_+)$, and define the equivalent relation $\mathcal{J}_+$ on V_0 by

$$x\mathcal{J}_+y \Longleftrightarrow x \text{ and } y \text{ belong to the same connected component of } G_+.$$

For $1 \leq i \leq n+1$, let $I_{i,+}$ be the equivalent class of $\mathcal{J}_+$ that contains $R_\lambda(c_{\kappa(i)})$.

(b). Define $\breve{G}_- = (C \cup R_\lambda(C), \breve{E}_-)$, with the edge set

$$\breve{E}_- = \{\{c_{i-1}, R_\lambda(c_{\kappa(i)})\} : 1 \leq i \leq n+1\}.$$

Define $G_- = \bigcup_{k=1}^\infty R^k_\lambda(\breve{G}_-)$, and define the equivalent relation $\mathcal{J}_-$ on V_0 by

$$x\mathcal{J}_-y \Longleftrightarrow x \text{ and } y \text{ belong to the same connected component of } G_-.$$

For $1 \leq i \leq n+1$, let $I_{i,-}$ be the equivalent class of $\mathcal{J}_-$ that contains $R_\lambda(c_{\kappa(i)})$.

Remark. Since κ is a permutation on $\{1, 2, \cdots, n+1\}$, its inversion, denoted as κ^{-1}, is well-defined. In particular, we have $R_\lambda(c_i) \in F_{\kappa^{-1}(i)}K_\lambda$, and $I_{\kappa^{-1}(i),+}$ is the equivalent class of $\mathcal{J}_+$ containing $R_\lambda(c_i)$. Similarly, $I_{\kappa^{-1}(i),-}$ is the equivalent class of $\mathcal{J}_-$ containing $R_\lambda(c_i)$.

We shall see that $\mathcal{J}_+$ and $\mathcal{J}_-$ are the only possibilities for non-trivial preserved $\mathcal{G}$-relations.

Lemma 3.20. *Let $\mathcal{J}$ be a non-trivial preserved $\mathcal{G}$-relation on V_0, then we have either $\mathcal{J} = \mathcal{J}_+$ or $\mathcal{J} = \mathcal{J}_-$. In addition, we always have $n+1$ disjoint equivalent classes I_i such that*

$$R_\lambda(c_{\kappa(i)}) \in I_i, \quad V_0 = \bigcup_{i=1}^{n+1} I_i.$$

Proof. Let $\mathcal{J}$ be a non-trivial preserved $\mathcal{G}$-relation on V_0. According to Lemma 3.11 (c), Lemma 3.18, and Lemma 3.12, there are $c_i \in C$ and $x \in R_\lambda(C)$ such that $c_i\mathcal{J}^{(1)}x$. In addition, by Lemma 3.12 and Lemma 3.11 (c), we have $x \in F_iV_0$ or $x \in F_{i+1}V_0$. In other words, either $c_i\mathcal{J}^{(1)}R_\lambda(c_{\kappa(i)})$ or $c_i\mathcal{J}^{(1)}R_\lambda(c_{\kappa(i+1)})$. By the rotation symmetry, we can see that either $\breve{E}_+ \subset \mathcal{J}^{(1)}$ or $\breve{E}_- \subset \mathcal{J}^{(1)}$, so by Lemma 3.10 (b),

$$\mathcal{J}_+ \subset \mathcal{J} \text{ or } \mathcal{J}_- \subset \mathcal{J}.$$

Without loss of generality, we may assume $\mathcal{J}_+ \subset \mathcal{J}$, and we need to show that $\mathcal{J}_+ = \mathcal{J}$. For any $x \in V_0$, there exists $1 \leq i \leq n+1$ and $k \geq 1$ such that $R_\lambda^k(c_{\kappa(i)}) = x$. Then, we have

$$x = R_\lambda^k(c_{\kappa(i)})\mathcal{J}_+R_\lambda^{k-1}(c_i)\mathcal{J}_+R_\lambda^{k-2}(c_{\kappa^{-1}(i)})\mathcal{J}_+\cdots\mathcal{J}_+R_\lambda(c_{\kappa^{2-k}(i)}),$$

which implies that $x \in I_{\kappa^{1-k}(i),+}$. So

$$V_0 = \bigcup_{i=1}^{n+1} I_{i,+}.$$

If $\mathcal{J}_+ \neq \mathcal{J}$, then there exists an equivalent class I of $\mathcal{J}$ and $i \neq j$ such that

$$I_{i,+} \cup I_{j,+} \subset I.$$

This implies that $R_\lambda(c_{\kappa(i)})\mathcal{J}R_\lambda(c_{\kappa(j)})$, and thus by Lemma 3.12, $j = i+1$ (without loss of generality) and $R_\lambda(c_{\kappa(i)})\mathcal{J}^{(1)}c_i\mathcal{J}^{(1)}R_\lambda(c_{\kappa(i+1)})$. By rotation symmetry, we also have $R_\lambda(c_{\kappa(i+1)})\mathcal{J}^{(1)}c_{i+1}$. Thus, $c_i\mathcal{J}^{(1)}c_{i+1}$, which is a contradiction by Lemma 3.11 (c). So we have $\mathcal{J}_+ = \mathcal{J}$.

Lastly, we claim that the equivalent classes $I_{i,+}, 1 \leq i \leq n+1$ are disjoint. Otherwise, by a same argument as above, we can see that $c_i\mathcal{J}^{(1)}c_{i+1}$ for some i. □

Next, we roughly describe the structure of $\mathcal{J}^{(1)}$.

Lemma 3.21. *Let $\mathcal{J}$ be a non-trivial preserved $\mathcal{G}$-relation, and I_i's be the equivalent classes of $\mathcal{J}$ as in Lemma 3.20.*

(a). For $(i,i') \neq (j,j') \in \{1,\cdots,n+1\}^2$, we have

$$x\mathcal{J}^{(1)}y, \forall x \in F_iI_{i'}, y \in F_jI_{j'} \Longleftrightarrow j = i+1, i' = j' = \kappa^{-1}(i) \text{ or } i = j+1, i' = j' = \kappa^{-1}(j).$$

In particular, we have that $\mathcal{J}^{(1)}$ consists of $n(n+1)$ equivalent classes.

(b). Define $I_i^{(1)}$ to be the equivalent class of $\mathcal{J}^{(1)}$ that contains I_i. We have

$$I_i^{(1)} = \begin{cases} F_iI_{\kappa^{-1}(i)} \cup F_{i+1}I_{\kappa^{-1}(i)}, & \text{if } \mathcal{J} = \mathcal{J}_+, \\ F_iI_{\kappa^{-1}(i)} \cup F_{i-1}I_{\kappa^{-1}(i)}, & \text{if } \mathcal{J} = \mathcal{J}_-. \end{cases}$$

In addition, $c_i \in I_i^{(1)}$ for each $i \in \{1,\cdots,n+1\}$.

Proof. (a). '$\Longleftarrow$': Let $1 \leq i \leq n+1$. Noticing that $R_\lambda(c_i) \in I_{\kappa^{-1}(i)}$ and $c_i = F_i(R_\lambda(c_i)) = F_{i+1}(R_\lambda(c_i))$, we have $c_i \in F_iI_{\kappa^{-1}(i)} \cap F_{i+1}I_{\kappa^{-1}(i)}$, so $F_iI_{\kappa^{-1}(i)}$ and $F_{i+1}I_{\kappa^{-1}(i)}$ are subsets of a same $\mathcal{J}^{(1)}$ class.
'$\Longrightarrow$': First, let's show that $i' = j'$. In fact, if the left side holds, then $R_\lambda(x)\mathcal{J}R_\lambda(y), \forall x \in F_iI_{i'}, y \in F_jI_{j'}$ by Lemma 3.10 (b), which is only possible if $i' = j'$, noticing that $R_\lambda(x) \in I_{i'}$ and $R_\lambda(y) \in I_{j'}$.

Next, we apply Lemma 3.12 to see that either $i = j$ or $|i - j| = 1$. However, the former case is impossible since it would imply $(i, i') = (j, j')$. Thus, $j = i + 1$ or $i = j + 1$. Without loss of generality, we assume $j = i + 1$. Then, by Lemma 3.12, we have $x\mathcal{J}^{(1)}c_i\mathcal{J}^{(1)}y$ for $x \in F_iI_{i'}, y \in F_jI_{i'}$. Thus $R_\lambda(x)\mathcal{J}R_\lambda(c_i)$, which implies that $i' = \kappa^{-1}(i)$ since $R_\lambda(x) \in I_{i'}$ and $R_\lambda(c_i) \in I_{\kappa^{-1}(i)}$.

From this conclusion it is easy to see that there are $(n+1)^2 - (n+1)$ equivalent classes of $\mathcal{J}^{(1)}$, as there are $n+1$ pairs of sets of the form $F_iI_{i'}$ matched.

(b). We look at the case $\mathcal{J} = \mathcal{J}_+$ only, since the argument for $\mathcal{J} = \mathcal{J}_-$ is the same. In this case, we have $c_i\mathcal{J}^{(1)}R_\lambda(c_{\kappa(i)})$ by definition, so $c_i \in I_i^{(1)}$, which implies

$$F_iI_{\kappa^{-1}(i)} \cup F_{i+1}I_{\kappa^{-1}(i)} \subset I_i^{(1)}.$$

Clearly, the left side itself is an equivalent class of $\mathcal{J}^{(1)}$ by (a). □

Now we can prove the existence of a $\mathcal{G}$-symmetric form on K_λ when $n \geq 2$, $m = 1$.

Proposition 3.22. *Let $R_\lambda(z) = z^n + \frac{\lambda}{z}$ with $n \geq 2$ be a MS map. Let $\mathcal{J}$ be a non-trivial preserved $\mathcal{G}$-relation on V_0. Then we have*

(a). $\overline{\rho}_{\mathcal{J}}^{\mathcal{G}} \leq 1$.

(b). $\underline{\rho}_{V_0/\mathcal{J}}^{\mathcal{G}} \geq 1 + \frac{1}{n} > 1$.

In particular, there exists exactly one $\mathcal{G}$-symmetric solution to (3.2).

Proof. (a). We abbreviate $x_i = R_\lambda(c_{\kappa(i)})$ below. Let I_i be the equivalent class of $\mathcal{J}$ containing x_i as in Lemma 3.20. Define a form $\mathcal{D} \in \mathcal{M}_{\mathcal{J}}$ as

$$\mathcal{D}(f) = \sum_{i=1}^{n+1} \sum_{x \in I_i \setminus \{x_i\}} \big(f(x) - f(x_i)\big)^2, \quad \forall f \in l(V_0).$$

Now, for each $f \in l(V_0)$, we define the extension $f_1 \in l(V_1)$ of f to be

$$f_1(x) = \begin{cases} f(x), & \text{if } x \in V_0, \\ f(x_i), & \text{if } x \in I_i^{(1)} \setminus I_i, \\ 0, & \text{if } x \in V_1 \setminus \bigcup_{i=1}^{n+1} I_i^{(1)}. \end{cases}$$

If $\mathcal{J} = \mathcal{J}_+$, by using Lemma 3.21, we have

$$\begin{aligned} \mathcal{D}^{(1)}(f_1) &= \sum_{i=1}^{n+1}\sum_{j=1}^{n+1} \sum_{x \in I_j \setminus \{x_j\}} \big(f_1(F_ix) - f_1(F_ix_j)\big)^2 \\ &= \sum_{i=1}^{n+1} \sum_{x \in I_{\kappa^{-1}(i)}} \big(f_1(F_ix) - f_1(c_i)\big)^2 + \sum_{i=1}^{n+1} \sum_{x \in I_{\kappa^{-1}(i-1)}} \big(f_1(F_ix) - f_1(c_{i-1})\big)^2 \\ &= \sum_{i=1}^{n+1} \sum_{x \in I_i^{(1)} \setminus \{c_i\}} \big(f_1(x) - f_1(c_i)\big)^2 = \sum_{i=1}^{n+1} \sum_{x \in I_i \setminus \{x_i\}} \big(f(x) - f(x_i)\big)^2 = \mathcal{D}(f). \end{aligned}$$

A similar equation holds for the case $\mathcal{J} = \mathcal{J}_-$. Using the above equation we have

$$\overline{\rho}_{\mathcal{J}}^{\mathcal{G}} \leq \overline{\rho}_{\mathcal{J}}^{\mathcal{G}}(\mathcal{D}) = \sup_{f \in l(V_0) \setminus l(V_0/\mathcal{J})} \frac{T_{\mathcal{J}}\mathcal{D}(f)}{\mathcal{D}(f)} \leq \sup_{f \in l(V_0) \setminus l(V_0/\mathcal{J})} \frac{\mathcal{D}^{(1)}(f_1)}{\mathcal{D}(f)} = 1.$$

(b). We define a form $\mathcal{D}\in\mathcal{M}_{V_0/\mathcal{J}}$ as

$$\mathcal{D}(f)=\sum_{i=1}^{n+1}\big(f(I_{i+1})-f(I_i)\big)^2,\quad \forall f\in l(V_0/\mathcal{J}).$$

Recall that we identify $l(V_0/\mathcal{J})$ with the subspace of $l(V_0)$ consisting of functions with constant value on each I_i, and identify $l(V_1/\mathcal{J}^{(1)})$ with a subspace of $l(V_1)$ analogously. Let f_1 be an extension of f to $l(V_1/\mathcal{J}^{(1)})$, we have

$$\begin{aligned}\mathcal{D}^{(1)}(f_1)&=\sum_{i=1}^{n+1}\sum_{j=1}^{n+1}\big(f_1(F_iI_{j+1})-f_1(F_iI_j)\big)^2\\&=\sum_{i=1}^{n+1}\Big(f_1\big(F_iI_{\kappa^{-1}(i-1)}\big)-f_1\big(F_iI_{\kappa^{-1}(i)}\big)\Big)^2+\sum_{i=1}^{n+1}\sum_{j\neq\kappa^{-1}(i-1)-1}\Big(f_1\big(F_iI_{j+1}\big)-f_1\big(F_iI_j\big)\Big)^2\\&=\sum_{i=1}^{n+1}\big(f(I_{i+1})-f(I_i)\big)^2+\sum_{i=1}^{n+1}\sum_{j\neq\kappa^{-1}(i)}\Big(f_1\big(F_iI_{j+1}\big)-f_1\big(F_iI_j\big)\Big)^2,\end{aligned}$$

where we use the property that $\kappa^{-1}(i-1)-1=\kappa^{-1}(i)$ (in cyclic notation $n+1=0$, follows from using Lemma 2.1 (a)) in the second equality, and Lemma 3.21 (b) in the last equality. By an effective resistance computation with series connection for the second term in the last line, also by Lemma 2.7, we get

$$\mathcal{D}^{(1)}(f_1)\geq(1+\frac{1}{n})\sum_{i=1}^{n+1}\big(f(I_{i+1})-f(I_i)\big)^2=(1+\frac{1}{n})\mathcal{D}(f).$$

Thus, we have

$$\underline{\rho}^{\mathcal{G}}_{V_0/\mathcal{J}}\geq\underline{\rho}^{\mathcal{G}}_{V_0/\mathcal{J}}(\mathcal{D})=\inf_{f\in l(V_0/\mathcal{J})}\frac{T_{\mathcal{J}}\mathcal{D}(f)}{\mathcal{D}(f)}\geq 1+\frac{1}{n}.$$

Finally, the existence and uniqueness of a $\mathcal{G}$-symmetric solution of (3.2) follows from Sabot's theorem, Theorem 3.7 (b). □

3.3. **Proof of uniqueness.** In this subsection, we return to the general MS Julia sets K_λ and prove that the $\mathcal{G}$-symmetric solution of (3.2), which has been shown to exist, is the unique solution (without assuming $\mathcal{G}$-symmetry a priori). For this aim, we will consider the non-trivial preserved relations $\mathcal{J}$ on V_0, which are not assumed to be $\mathcal{G}$-symmetric. Luckily, we can take the advantage of the existence of a symmetric form and the 'ring' structure of the level-1 cells.

Throughout this subsection, we would admit the following settings:

We fix a solution $\mathcal{D}$ to (3.2) which has already been proved to exist. In particular, we have a positive constant η (which is uniquely determined by the equation (3.2)), such that

$$T\mathcal{D}=\eta^{-1}\mathcal{D}.$$

Note that $0<\eta^{-1}<1$ by a well-known theorem (see [21, 22]), and the pair $(\mathcal{D},\eta^{-1})$ is called a *regular harmonic structure* on K_λ, which will generate a local resistance form $(\mathcal{E},\mathcal{F})$ on K_λ. For reference, a standard proof is available in Kigami's book [22].

We will show that for any preserved relation on V_0,

$$\overline{\rho}_{\mathcal{J},k} < \underline{\rho}_{V_0/\mathcal{J},k}$$

for some $k \geq 1$, where $\overline{\rho}_{\mathcal{J},k}$ and $\underline{\rho}_{V_0/\mathcal{J},k}$ are defined in the remark at the end of Subsection 3.1. In particular, we will construct forms based on the solution $\mathcal{D}$.

Recalling the definition of $\mathcal{D}_{V_0/\mathcal{J}}$, the following easy lemma follows from Sabot's paper [35].

Lemma 3.23. *Let $\mathcal{J}$ be a non-trivial preserved relation on V_0, and let $f \in l(V_0/\mathcal{J})$. We have*

$$T_{V_0/\mathcal{J}}\mathcal{D}_{V_0/\mathcal{J}}(f) \geq T\mathcal{D}(f) = \eta^{-1}\mathcal{D}(f).$$

As an easy consequence, we can see

Lemma 3.24. *Let $\mathcal{J}$ be a non-trivial preserved relation on V_0, then $\underline{\rho}_{V_0/\mathcal{J},k} \geq \eta^{-k}$ for any $k \geq 1$.*

We will devote the rest of this subsection to demonstrate a corresponding statement:

Proposition 3.25. *Let $\mathcal{J}$ be a non-trivial preserved relation on V_0, then $\overline{\rho}_{\mathcal{J},k} < \eta^{-k}$ for some $k \geq 1$.*

To prove the proposition, we will study the following form in $\mathcal{M}_{\mathcal{J}}$.

Definition 3.26. *Let $\mathcal{J}$ be a preserved relation on V_0, with equivalent classes I_a, $a = 1, 2, \cdots, N$.*

(a). We define

$$\mathcal{D}_{\mathcal{J}} = \sum_{a=1}^{N} \mathcal{D}|_{I_a}.$$

(b). For $1 \leq a \leq N$, we define

$$\ell(I_a) = \{f \in l(V_0) : f|_{V_0 \setminus I_a} \equiv 0\}.$$

(c). For $1 \leq a \leq N$ and $f \in \ell(I_a)$, we define $H^{(1)}_{\mathcal{D}_{\mathcal{J}},I_a} f$ to be the unique function in $l(V_1)$ such that

$$(H^{(1)}_{\mathcal{D}_{\mathcal{J}},I_a} f)|_{V_1 \setminus I_a^{(1)}} \equiv 0,$$

and

$$\mathcal{D}^{(1)}_{\mathcal{J}}\big(H^{(1)}_{\mathcal{D}_{\mathcal{J}},I_a} f\big) = T_{\mathcal{J}}\mathcal{D}_{\mathcal{J}}(f),$$

where we recall that $\mathcal{D}^{(1)}_{\mathcal{J}}(g) = \sum_{i=1}^{m+n} \mathcal{D}_{\mathcal{J}}(g \circ F_i)$ for any $g \in l(V_1)$.

Since we no longer are assuming rotational symmetry, our starting points will be Lemmas 3.10 and 3.11. The following statement is an easy consequence of Lemma 3.10.

Lemma 3.27. *Let $\mathcal{J}$ be a non-trivial preserved relation on V_0, with equivalent classes $I_a, a = 1, 2, \cdots, N$. For any $1 \leq a \leq N$ and $k \geq 1$, there is a unique $\mathcal{J}^{(k)}$ class $I_a^{(k)}$ such that $I_a \subset I_a^{(k)}$. In addition, the classes $I_a^{(k)}, a = 1, 2, \cdots, N$ are disjoint. Finally, for any finite word w with $|w| = k$, if $F_w V_0 \cap I_a^{(k)} \neq \emptyset$, we have a unique $a_w \in \{1, 2, \cdots, N\}$ such that $F_w I_{a_w} \subset I_a^{(k)}$.*

Proof. The existence of $I_a^{(k)}$ and the claim that $I_a^{(k)}$ are disjoint follow quickly from Lemma 3.10 (a). To see the last statement, we assume there exist $I_b \neq I_{b'}$ such that $F_w I_b \subset I_a^{(k)}, F_w I_{b'} \subset I_a^{(k)}$. Then for $x \in I_b$ and $y \in I_{b'}$, using Lemma 3.10 (b), we have

$$F_w x \mathcal{J}^{(k)} F_w y \Longrightarrow x \mathcal{J} y.$$

This is impossible because x and y have been chosen from different equivalent classes of $\mathcal{J}$. □

The above lemma allows us to prove the following result.

Lemma 3.28. *Let $\mathcal{J}$ be a preserved relation on V_0, with equivalent classes $I_a, a = 1, 2, \cdots, N$.*

(a). Let $a \neq b$. Then

$$\tilde{\mathcal{D}}(f, g) = 0, \quad \forall \tilde{\mathcal{D}} \in \mathcal{M}_{\mathcal{J}}, f \in \ell(I_a), g \in \ell(I_b).$$

(b). Let $k \geq 1$ and $f \in l(V_0)$, we have

$$\eta^k T_{\mathcal{J}}^k \mathcal{D}_{\mathcal{J}}(f) \leq \eta^{k-1} T_{\mathcal{J}}^{k-1} \mathcal{D}_{\mathcal{J}}(f) \leq \cdots \leq \mathcal{D}_{\mathcal{J}}(f). \tag{3.5}$$

Moreover, for fixed $1 \leq a \leq N$, let $h_f := H_{\mathcal{E}, I_a}(f) \in C(K_\lambda)$, which is the harmonic extension of $f|_{I_a}$ with respect to the form $(\mathcal{E}, \mathcal{F})$ generated by $\mathcal{D}$. Also, let $W_{k,a} = \{|w| = k : F_w V_0 \cap I_a^{(k)} \neq \emptyset\}$. Then,

$$\eta^k T_{\mathcal{J}}^k \mathcal{D}_{\mathcal{J}}(f) = \mathcal{D}_{\mathcal{J}}(f) \tag{3.6}$$

only if

$$h_f \circ F_w = \begin{cases} H_{\mathcal{E}, I_{a_w}}\big((h_f \circ F_w) \cdot 1_{I_{a_w}}\big), & \text{if } w \in \bigcup_{l=1}^k W_{l,a}, \\ \text{constant}, & \text{if } 1 \leq |w| \leq k, w \notin \bigcup_{l=1}^k W_{l,a}. \end{cases} \tag{3.7}$$

where I_{a_w} is the unique $\mathcal{J}$ class such that $F_w I_{a_w} \subset I_a^{(|w|)}$ as shown in Lemma 3.27.

Proof. (a) is obvious. We will focus on (b). Fix $1 \leq a \leq N$ and $f \in \ell(I_a)$, then we can show the $k = 1$ case of (3.5) by the following computation,

$$\begin{aligned} \eta T_{\mathcal{J}} \mathcal{D}_{\mathcal{J}}(f) &= \eta \mathcal{D}_{\mathcal{J}}^{(1)}(H_{\mathcal{D}_{\mathcal{J}}, I_a}^{(1)} f) \\ &\leq \eta \mathcal{D}_{\mathcal{J}}^{(1)}(h_f \cdot 1_{I_a^{(1)}}) = \eta \sum_{i=1}^{m+n} \mathcal{D}_{\mathcal{J}}\big((h_f \cdot 1_{I_a^{(1)}}) \circ F_i\big) \\ &= \eta \sum_{i \in W_{1,a}} \mathcal{D}_{\mathcal{J}}\big((h_f \circ F_i) \cdot 1_{a_i}\big) \\ &\leq_{(*1)} \eta \sum_{i=1}^{m+n} \mathcal{D}(h_f \circ F_i) = \mathcal{D}(h_f|_{V_0}) = \mathcal{D}|_{I_a}(f) = \mathcal{D}_{\mathcal{J}}(f), \end{aligned} \tag{3.8}$$

where in the first inequality we use the fact that $H_{\mathcal{D}_{\mathcal{J}}, I_a}^{(1)} f$ is the minimal energy extension, and in the third equality and the second inequality we use Lemma 3.27. For a general $f \in l(V_0)$, we apply (a) to show (3.5) still holds,

$$\eta T_{\mathcal{J}} \mathcal{D}_{\mathcal{J}}(f) = \eta \sum_{a=1}^N T_{\mathcal{J}} \mathcal{D}_{\mathcal{J}}(f \cdot 1_{I_a}) \leq \sum_{a=1}^N \mathcal{D}_{\mathcal{J}}(f \cdot 1_{I_a}) = \mathcal{D}_{\mathcal{J}}(f),$$

since $\ell(I_a)$'s are pairwise orthogonal for different a's with respect to forms in $\mathcal{M}_{\mathcal{J}}$ by (a).

Next, we return to study the conditions under which (3.5) becomes an equality. Clearly, for (3.8) to be an equality, we need '$\leq_{(*1)}$' holds, which is equivalent to the condition (3.7) for $k=1$. To extend the observation to general $k \geq 1$, we induct.

Assuming that (3.6) holds, we have immediately that $\eta T_{\mathcal{J}}\mathcal{D}_{\mathcal{J}}(f) = \mathcal{D}_{\mathcal{J}}(f)$ by (3.5). So (3.7) holds immediately for $|w|=1$. Next, we apply the inductive assumption to get the following inequality,

$$\begin{aligned}
\eta^k T^k_{\mathcal{J}}\mathcal{D}_{\mathcal{J}}(f) &= \eta^k (T^{k-1}_{\mathcal{J}}\mathcal{D}_{\mathcal{J}})^{(1)}(H^{(1)}_{T^{k-1}_{\mathcal{J}}\mathcal{D}_{\mathcal{J}},I_a} f)\\
&\leq \eta^k (T^{k-1}_{\mathcal{J}}\mathcal{D}_{\mathcal{J}})^{(1)}(h_f \cdot 1_{I_a^{(1)}})\\
&= \eta \sum_{i\in W_{1,a}} \eta^{k-1} T^{k-1}_{\mathcal{J}}\mathcal{D}_{\mathcal{J}}\big((h_f\cdot 1_{I_a^{(1)}})\circ F_i\big)\\
&\leq_{(*2)} \eta \sum_{i\in W_{1,a}} \mathcal{D}_{\mathcal{J}}\big((h_f\cdot 1_{I_a^{(1)}})\circ F_i\big)\\
&= \eta \sum_{i\in W_{1,a}} \mathcal{D}\big(h_f\circ F_i\big) = \mathcal{D}(h_f|_{V_0}) = \mathcal{D}|_{I_a}(f) = \mathcal{D}_{\mathcal{J}}(f),
\end{aligned}$$

where we use the fact $h_f\circ F_i = H_{\mathcal{E},I_{a_i}}\big((h_f\circ F_i)\cdot 1_{I_{a_i}}\big) = H_{\mathcal{E},I_{a_i}}\big((h_f\cdot 1_{I_a^{(1)}})\circ F_i\big)$ by (3.7) in the last line.

Clearly, when (3.6) holds, we require '$\leq_{(*2)}$' to be an equality. This means

$$\eta^{k-1}T^{k-1}_{\mathcal{J}}\mathcal{D}_{\mathcal{J}}\big((h_f\cdot 1_{I_a^{(1)}})\circ F_i\big) = \mathcal{D}_{\mathcal{J}}\big((h_f\cdot 1_{I_a^{(1)}})\circ F_i\big), \quad \forall i\in W_{1,a}.$$

By the inductive hypothesis, this implies that (3.7) holds for $1\leq |w|\leq k-1$, with $h_f\circ F_i$ replacing h_f and a_i replacing a for $i\in W_{1,a}$, and thus (3.7) holds for any $2\leq |w|\leq k$. $\square$

To understand the condition (3.7) better, we introduce the concept of flows. The concept is also known as that of normal derivatives for its role in the Gauss-Green's formula, but we prefer the word 'flow' for the more intuitive physical picture.

Definition 3.29. *Let h be a harmonic function on K_λ, i.e. $\mathcal{E}(h) = \mathcal{D}(h|_{V_0})$.*

(a). For $x\in V_0$, we define the flow (normal derivative) of h at x as $\partial_n h(x) = \mathcal{D}(h, 1_x)$. We write $V_{0,h} := \{x\in V_0 : \partial_n h(x)\neq 0\}$ for the set of boundary vertices with nonzero flows.

(b). For $c_i\in C$, we say h has nonzero flow passing through c_i if

$$\partial_n(h\circ F_j)\big(R_\lambda(c_i)\big)\neq 0,\ \text{for } j=i,i+1.$$

We write $C_h := \{c_i\in C : \partial_n(h\circ F_i)\big(R_\lambda(c_i)\big)\neq 0\}$ for the set of critical points with nonzero flows.

We enumerate some simple properties of flow.

(P1). $\sum_{x\in V_0}\partial_n h(x) = 0$.

(P2). $\partial_n(h\circ F_i)\big(R_\lambda(c_i)\big) + \partial_n(h\circ F_{i+1})\big(R_\lambda(c_i)\big) = 0$.

(P3). For $x\in V_0\cap F_i(V_0)$, we have the scaling identity $\partial_n h(x) = \eta\partial_n(h\circ F_i)\big(R_\lambda(x)\big)$.

The following observation is obvious.

Lemma 3.30. *Let $\mathcal{J}$ be a non-trivial preserved relation on V_0, with equivalent classes $I_a, a = 1,2,\cdots,N$. Let h be a harmonic function on K_λ and $1\leq a\leq N$. If (3.7) holds for $k=1$ with h replacing h_f, we have $C_h\subset I_a^{(1)}$, which implies that $C_h\neq C$ by Lemma 3.11 (a).*

Next, we prove that $T^k_{\mathcal{J}}\mathcal{D}_{\mathcal{J}} < \mathcal{D}_{\mathcal{J}}$ for some $k \geq 1$ by means of Lemmas 3.28 and 3.30. First, we show an estimate for a single $f \in l(V_0)$.

Lemma 3.31. *Let $\mathcal{J}$ be a non-trivial preserved relation on V_0, and $f \in l(V_0) \setminus l(V_0/\mathcal{J})$. Then there exists $k \geq 1$ such that*

$$\eta^k T^k_{\mathcal{J}}\mathcal{D}_{\mathcal{J}}(f) < \mathcal{D}_{\mathcal{J}}(f).$$

Proof. We begin by taking $f \in \ell(I_a)$ for some $1 \leq a \leq N$ which is non-constant on I_a. We will prove the lemma by contradiction. Assume $\eta^k T^k_{\mathcal{J}}\mathcal{D}_{\mathcal{J}}(f) = \mathcal{D}_{\mathcal{J}}(f), \forall k \geq 0$. Then, by Lemmas 3.28 and 3.30, we have

$$\#C_{h_f \circ F_w} < m + n, \tag{3.9}$$

for any finite word w and $h_f = H_{\mathcal{E},I_a}(f)$. We will see this is impossible for some word w. We will construct such a word in two steps.

Step 1. *We can find a finite word w such that $\#V_{0,h_f\circ F_w} = 2$.*

To achieve this, we will construct a sequence of finite words w_k of length k inductively. For convenience, we write $N_k := \#V_{0,h_f\circ F_{w_k}}$ and $M_k := \#C_{h_f\circ F_{w_k}}$. We start from $w_0 = \emptyset$, and choose $w_1, w_2, \cdots$ in accordance with the following rules, stopping when $N_k = 2$.

Assume we have chosen w_k with $N_k \geq 3$, we will choose $w_{k+1} = w_k i$ with $1 \leq i \leq m+n$ following two possible cases.

Case 1.1: $M_k = 0$. In this case, we simply choose an $1 \leq i \leq m+n$ such that $\#V_{0,h_f\circ F_{w_k i}} > 0$. Let $w_{k+1} = w_k i$ and clearly we have $\#N_{k+1} \leq \#N_k$.

Case 1.2: $M_k \geq 1$. In this case, by (3.9), we also have $M_k \leq m+n-1$. So, there are at least $M_k + 1$ different i's such that $\#V_{0,h_f\circ F_{w_k i}} > 0$. Moreover,

$$\sum_{i=1}^{m+n} \#V_{0,h_f\circ F_{w_k i}} = 2M_k + N_k.$$

Thus, we can choose $w_{k+1} = w_k i$ such that

$$N_{k+1} = \#V_{0,h_f\circ F_{w_{k+1}}} \leq \frac{2M_k + N_k}{M_k + 1} < \frac{N_k M_k + N_k}{M_k + 1} = N_k, \tag{3.10}$$

where '$<$' holds since we always have $N_k \geq 3$ (otherwise, we will stop the construction) and $M_k \geq 1$.

Continuing the construction, we can easily see that $N_0 \geq N_1 \geq N_2 \geq \cdots$. However, Case 1.1 cannot repeat consecutively for infinitely many iterations, otherwise, since it does not introduce any new flow, after long iterations we will have a small cell $F_w K_\lambda$ which contains only 1 original flow at its boundary, which is impossible by (P1). Each time we face Case 1.2, we have a strict decrease in N_k. Eventually, we can find a $k \geq 1$ such that $\#N_k = 2$.

Step 2. *We can find a finite word w such that $V_{0,h_f\circ F_w} = \{x, y\}$, with $x \in F_i K_\lambda$, $y \in F_j K_\lambda$ and $i \neq j$.*

Let w be the word found in Step 1. If w satisfies the condition, we are done. Otherwise, we have $V_{0,h_f\circ F_w} \subset F_i K_\lambda$ for some $1 \leq i \leq m+n$. We may face either

Case 2.1: $C_{h_f\circ F_w} = \emptyset$, or *Case 2.2:* $C_{h_f\circ F_w} \neq \emptyset$.

Case 2.2 is clearly impossible since that would imply $C = C_{h_f\circ F_w}$ by the ring structure of 1-cells of K_λ. Hence, only Case 2.1 is possible, and we may choose $w' = wi$. After repeating the above argument finitely many times, we will find a finite word w'' satisfying the desired condition.

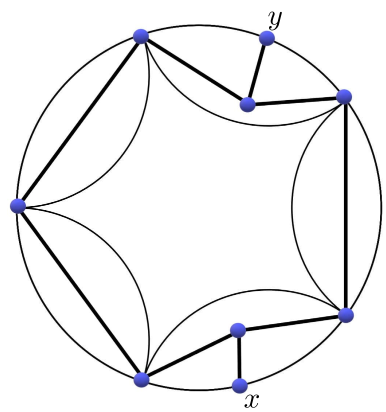


FIGURE 8. The restriction of $\mathcal{E}$ onto $C \cup \{x, y\}$.

Now, we look at the word w chosen in Step 2. There are $x, y \in V_0$ such that $h_f \circ F_w$ is harmonic in $K_\lambda \setminus \{x, y\}$. In particular, by restricting $\mathcal{E}$ to $\{x, y\} \cup C$ and applying the $\Delta - Y$ transformation (see books [22, 38] for the formulas of $\Delta - Y$ transformation) in the cells $F_i K_\lambda, F_j K_\lambda$ (see Figure 8 for an illustration of the restricted electrical network), we can easily see that $h_f \circ F_w$ has nonzero flow at each $c \in C$. This contradicts (3.9).

Finally, for a general $f \in l(V_0) \setminus l(V_0/\mathcal{J})$, there is at least one $1 \leq a \leq N$ such that $f \cdot 1_{I_a} \in \ell(I_a)$ is non-constant on I_a. So there is $k \geq 1$ such that $\eta^k T_{\mathcal{J}}^k \mathcal{D}_{\mathcal{J}}(f \cdot 1_{I_a}) < \mathcal{D}_{\mathcal{J}}(f \cdot 1_{I_a})$, and then,

$$\eta^k T_{\mathcal{J}}^k \mathcal{D}_{\mathcal{J}}(f) = \sum_{1 \leq a' \leq N} \eta^k T_{\mathcal{J}}^k \mathcal{D}_{\mathcal{J}}(f \cdot 1_{I_{a'}}) < \sum_{1 \leq a' \leq N} \mathcal{D}_{\mathcal{J}}(f \cdot 1_{I_{a'}}) = \mathcal{D}_{\mathcal{J}}(f),$$

which finishes the proof. □

We are now ready to prove Proposition 3.25.

Proof of Proposition 3.25. Define

$$B_{\mathcal{J}} = \{g \in l(V_0) : \max\{|g(x) - g(y)| : x\mathcal{J}y\} = 1, \text{ and } \|g\|_{l^\infty(V_0)} = 1/2\}$$

as a compact subset of $l^\infty(V_0)$. Then for any fixed $f \in l(V_0) \setminus l(V_0/\mathcal{J})$, we can find $g \in B_{\mathcal{J}} \cap \{cf + u : c \in \mathbb{R}, u \in l(V_0/\mathcal{J})\}$, and it is easy to see that

$$\frac{T_{\mathcal{J}}^k \mathcal{D}_{\mathcal{J}}(f)}{\mathcal{D}_{\mathcal{J}}(f)} = \frac{T_{\mathcal{J}}^k \mathcal{D}_{\mathcal{J}}(g)}{\mathcal{D}_{\mathcal{J}}(g)}, \quad \forall k \geq 1.$$

In addition, since $\mathcal{D}_{\mathcal{J}}(g) > 0$ on $B_{\mathcal{J}}$, the function $\frac{T_{\mathcal{J}}^k \mathcal{D}_{\mathcal{J}}(g)}{\mathcal{D}_{\mathcal{J}}(g)}$ is continuous on $B_{\mathcal{J}}$.

We now claim that $\eta^k T_{\mathcal{J}}^k \mathcal{D}_{\mathcal{J}} < \mathcal{D}_{\mathcal{J}}$ for some $k \geq 1$. Assume not, then for any $k \geq 1$, there exists $f_k \in l(V_0) \setminus l(V_0/\mathcal{J})$ such that $\eta^k T_{\mathcal{J}}^k \mathcal{D}_{\mathcal{J}}(f_k) = \mathcal{D}_{\mathcal{J}}(f_k)$. In addition, we can require that $f_k \in B_{\mathcal{J}}$ by the previous argument. Thus, there exists a subsequence $k_l, l \geq 1$ such that f_{k_l} converges to a function $f \in B_{\mathcal{J}}$. Clearly,

$$\eta^k T_{\mathcal{J}}^k \mathcal{D}_{\mathcal{J}}(f) = \lim_{l \to \infty} \eta^k T_{\mathcal{J}}^k \mathcal{D}_{\mathcal{J}}(f_{k_l}) = \lim_{l \to \infty} \mathcal{D}_{\mathcal{J}}(f_{k_l}) = \mathcal{D}(f), \quad \forall k \geq 1.$$

This contradicts Lemma 3.31.

Lastly, the proposition follows from the inequality

$$\overline{\rho}_{\mathcal{J},k} \leq \overline{\rho}_{\mathcal{J},k}(\mathcal{D}) = \sup_{f\in l(V_0)\setminus l(V_0/\mathcal{J})} \frac{T_{\mathcal{J}}^k \mathcal{D}_{\mathcal{J}}(f)}{\mathcal{D}_{\mathcal{J}}(f)} = \sup_{f\in B_{\mathcal{J}}} \frac{T_{\mathcal{J}}^k \mathcal{D}_{\mathcal{J}}(f)}{\mathcal{D}_{\mathcal{J}}(f)} < \eta^{-k}.$$

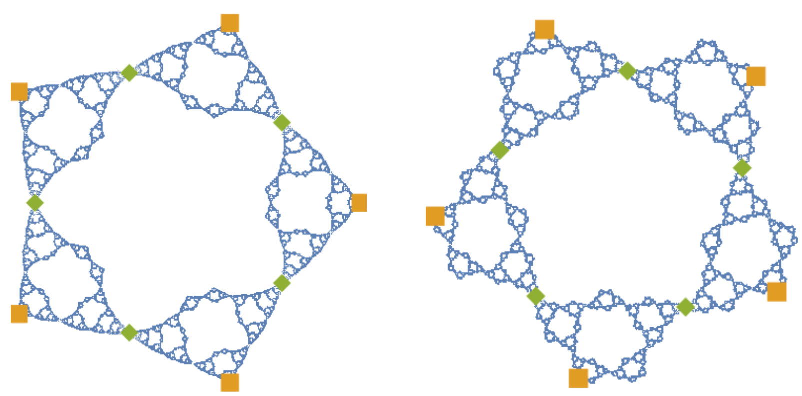
□

Finally, we conclude the proof of the main result in this section.

Proof of Theorem 3.2. The existence of a symmetric form follows from Proposition 3.16 and Proposition 3.22. The uniqueness follows from Lemma 3.24, Proposition 3.25, Theorem 3.7 and the remark after Theorem 3.7. □

3.4. **Examples.** The MS Julia sets can be quite complicated in general. Although we have established the existence and uniqueness theorem for the balanced resistance forms, we are only able to compute the exact forms for some simple examples.

Example 3.32. *Consider* $R_\lambda(z) = z^n + \frac{\lambda}{z^m}$ *with* $\theta_\lambda = \frac{l}{n(m+n)}$ *for* $m \geq 1$, $n \geq 2$ *and* $1 \leq l \leq n-1$. *For* m, n *fixed, these parameters correspond to the Julia sets with the smallest possible* $\tilde{V}_0 = \{\psi_\lambda^{-1}[\frac{k}{m+n}] : 1 \leq k \leq m+n-1\}$. *See Figure 9 for examples of such sets. These Julia sets turn out to be exactly the* N*-gaskets studied in* [8].



$m = 3, n = 2$ and $\theta_\lambda = \frac{1}{10}$ $m = 2, n = 3$ and $\theta_\lambda = \frac{2}{15}$

FIGURE 9. Some MS Julia sets with smallest $\tilde{V}_0$.

Due to Theorem 3.2, there exists exactly one balanced form on K_λ. *To simplify the calculation, we consider* $V_0' := \{p_0, p_1, p_2\}$ *with*

$$\psi_\lambda(p_0) = [0], \quad \psi_\lambda(p_1) = [\frac{l}{m+n}], \quad \psi_\lambda(p_2) = [\frac{m+l}{m+n}],$$

and set $V_1' = C \cup \tilde{V}_0$. *Note that* p_1, p_2 *are the two nearest vertices in* $\tilde{V}_0$ *surround* p_0. *By a simple computation (using the* $\Delta - Y$ *transformation to restrict the form on* V_1' *to* V_0'*), we get the exact value of the renormalization constant* η *to be*

$$\eta = \frac{1}{2} + \frac{mn}{2(m+n)} + \frac{1}{2}\sqrt{(\frac{mn}{m+n} - 1)^2 + \frac{8l(n-l)}{m+n}},$$

with the form $\mathcal{D}|_{V_0'}$ as shown in Figure 10. We omit the computation here, and readers can find a similar computation for the pentagasket in the book [38].

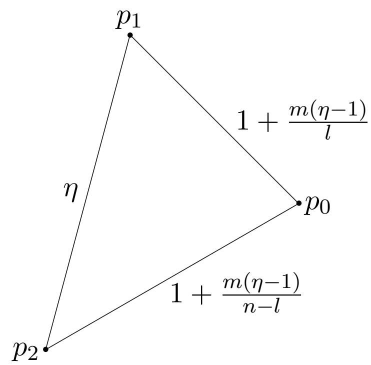


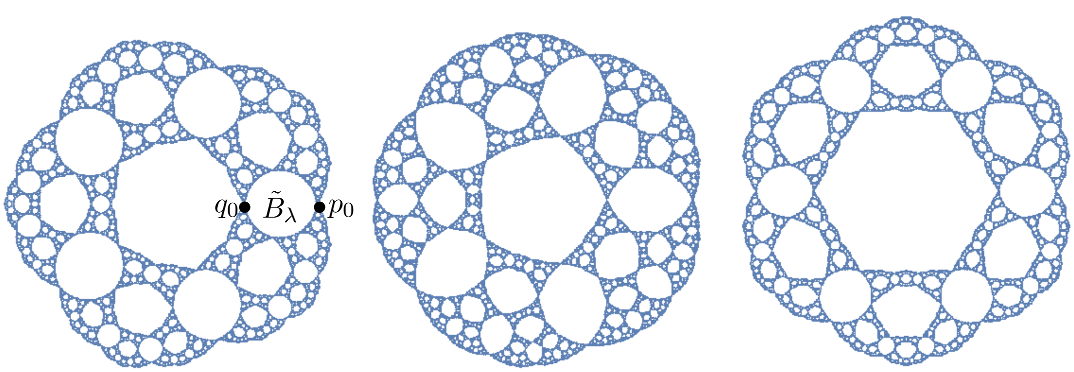
FIGURE 10. The restriction of $\mathcal{D}$ to V_0'.

Example 3.33. *For the Julia set in Example 2.6 (a), we have experimentally that $\eta^{-1} \approx 0.64735$. On the other hand, one can check easily that $T_{\mathcal{J}}\mathcal{D}_{\mathcal{J}} = \frac{1}{2}\mathcal{D}_{\mathcal{J}}$ for the relation $\mathcal{J}$ shown in Example 3.17. In particular, this shows $\eta T_{\mathcal{J}}\mathcal{D}_{\mathcal{J}} < \mathcal{D}_{\mathcal{J}}$ as Proposition 3.25 states.*

4. Other finitely ramified Julia sets of rational maps

In this last section of the paper, we look at some other Julia sets also associated to rational maps $R_\lambda(z) = z^n + \frac{\lambda}{z^m}$ with $n \geq 2$, $m \geq 1$, which are not MS maps. Instead of providing a full story as done in Section 3, this section is more explorative, and we hope that the observations may lead to further studies.

In particular, we focus on the simple class of rational maps R_λ whose critical set possesses a real fixed point c. By an easy calculation from $R'_\lambda(c) = 0$ and $R_\lambda(c) = c$, we have $\lambda = \frac{nc^{n+m}}{m}$ and $c = (\frac{n}{m+n})^{\frac{1}{n-1}}$ in this case. Clearly, c is a superattracting fixed point, so the immediate attracting basin of c is excluded from the Julia set K_λ, see [27, Section 9]. See Figure 11 for some examples of these Julia sets.



$m = 2, n = 3$ $\qquad$ $m = 3, n = 2$ $\qquad$ $m = 3, n = 3$

FIGURE 11. Julia sets with a fixed real critical point.

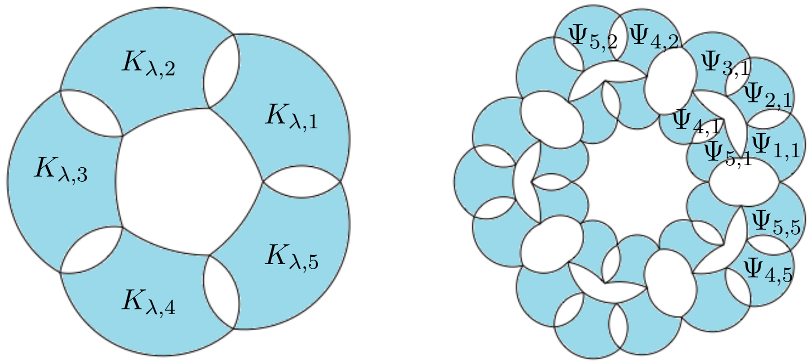


FIGURE 12. An illustration of level-1 cells and the $\Psi_{k,l}$ mappings for $(m, n) = (2, 3)$.

We let B_λ denote the immediate attracting basin of ∞, $T_\lambda = R_\lambda^{-1}(B_\lambda)$ denote the Fatou component containing 0, and $\tilde{B}_\lambda$ denote the immediate basin of the real c. Then we have two local cut points p_0, q_0, namely

$$\{p_0\} = \overline{B}_\lambda \cap \overline{\tilde{B}_\lambda}, \quad \{q_0\} = \overline{T}_\lambda \cap \overline{\tilde{B}_\lambda}.$$

Since the Julia set K_λ admits the rotation symmetry, we define

$$p_l = e^{\frac{2l\pi i}{m+n}} p_0, \; q_l = e^{\frac{2l\pi i}{m+n}} p_0, \text{ for } 0 \le l \le m+n-1.$$

The vertex set $\{p_l, q_l\}_{l=0}^{m+n-1}$ cuts the Julia set K_λ into $m+n$ connected components. We denote by $K_{\lambda,l}$ the closure of one of the components such that $K_{\lambda,l}$ contains $\{p_l, q_l, p_{l-1}, q_{l-1}\}$ for $l = 1, 2, \ldots m+n$, using the cyclic notation $m+n = 0$, and call them the 1-*cells* of K_λ.

It is not hard to see that $R_\lambda^{-1}(\{p_l, q_l\}_{l=0}^{m+n-1})$ is a set of $2(m+n)^2$ vertices that contains $\{p_l, q_l\}_{l=0}^{m+n-1}$ and cuts K_λ into $(m+n)^2$ pieces. For each pair of $1 \le k, l \le m+n$, we denote $\Psi_{k,l}$ for the local inverse of R_λ such that $\Psi_{k,l} : K_{\lambda,k} \to K_{\lambda,l}$. Then we have

$$K_{\lambda,l} = \bigcup_{k=1}^{m+n} \Psi_{k,l}(K_{\lambda,k}).$$

We still have the same conjugacy of R_λ on the boundary β_λ of B_λ to the angle mapping Φ_n on the unit circle $\mathbb{T}$ as described in Section 2, and we have $R_\lambda(p_l) = R_\lambda(q_l)$. These facts determine the position of each level-2 cells $\Psi_{k,l}(K_{\lambda,k})$. See Figure 12 for an illustration (where we take $m = 2, n = 3$).

By iterating the mappings in the proper way, we see that the diameters of the higher level cells shrink to 0. More precisely, we have

$$\mathrm{diam}(\Psi_{l_{k-1},l_k} \cdots \Psi_{l_1,l_2} \Psi_{l_0,l_1}(K_{\lambda,l_0})) \to 0, \text{ as } k \to \infty,$$

for any infinite sequence $\{l_k\}_{k\ge0}$. This provides us with a graph-directed structure of K_λ (see [9, 18]).

We consider the forms $\mathcal{E}_k, k = 1, 2, \cdots, m+n$ on $K_{\lambda,k}$'s satisfying the graph-directed invariance, i.e.

$$\mathcal{E}_k(f) = \eta \sum_{l=1}^{m+n} \mathcal{E}_l(f \circ \Psi_{l,k}), \tag{4.1}$$

for some positive constant η independent of k.

For simplicity, we still consider the rotationally symmetric solutions, i.e. we require

$$\mathcal{E}_k(f) = \mathcal{E}_{k+l}(f(e^{-\frac{2l\pi i}{m+n}}\bullet)) \tag{4.2}$$

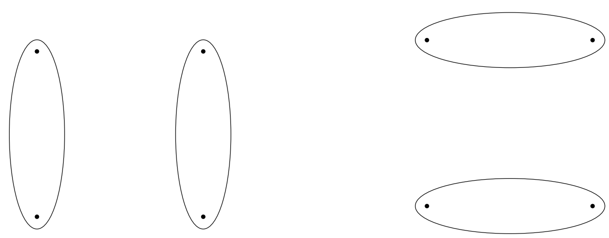
holds for any pair of $1 \leq k, l \leq m+n$. Then (4.1) is simplified to an equation of the form (3.2) (with different contractive mappings of course),

$$\mathcal{E}_1(f) = \eta \sum_{l=1}^{m+n} \mathcal{E}_1(f \circ \Psi_{l,1}(e^{\frac{2(l-1)\pi i}{m+n}}\bullet)). \tag{4.3}$$

In particular, the existence of a solution to (4.1) is equivalent to the existence of a rotationally symmetric solution to (4.1). This can be easily proven with Hilbert's projective metric [29] and the Brouwer fixed point theorem. See [3], Proposition 6.21, for a similar result on p.c.f. self-similar sets, whose proof can be easily modified for our purpose.

We again apply Sabot's criteria, Theorem 3.7, to study the existence of forms that satisfy (4.3). In particular, there are only two non-trivial preserved relations $\mathcal{J}_1$ and $\mathcal{J}_2$ on $V := \{p_0, p_1, q_0, q_1\}$, depicted in Figure 13.

1. $\mathcal{J}_1$ consists of two equivalent classes $\{p_0, q_0\}$ and $\{p_1, q_1\}$.
2. $\mathcal{J}_2$ consists of two equivalent classes $\{p_0, p_1\}$ and $\{q_0, q_1\}$.



$\mathcal{J}_1$ $\mathcal{J}_2$

FIGURE 13. The non-trivial preserved relations on $\{p_0, p_1, q_0, q_1\}$.

It is easy to compute the exact values of the $\overline{\rho}, \underline{\rho}$'s for $\mathcal{J}_1$ and $\mathcal{J}_2$.

Proposition 4.1. *For $\mathcal{J}_1, \mathcal{J}_2$ defined above, we have*

$$\underline{\rho}_{\mathcal{J}_1} = \overline{\rho}_{\mathcal{J}_1} = \frac{1}{2}, \quad \underline{\rho}_{V/\mathcal{J}_1} = \overline{\rho}_{V/\mathcal{J}_1} = \frac{1}{m} + \frac{1}{n}, \quad \overline{\rho}_{\mathcal{J}_2} = \frac{1}{n}, \quad \underline{\rho}_{V/\mathcal{J}_2} = \frac{mn}{m+n}.$$

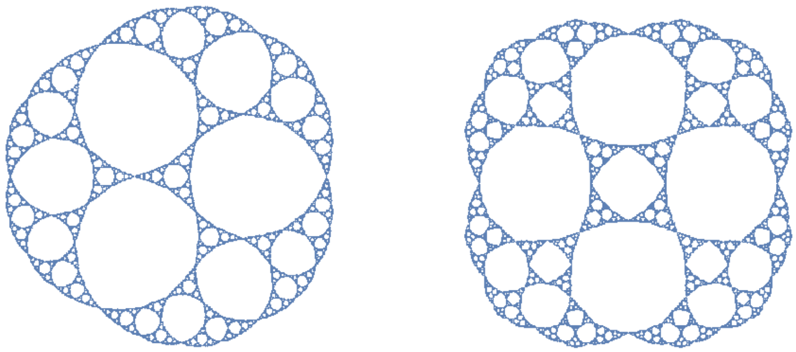
As a consequence, we have

(a). There are no graph-directed invariant forms on K_λ if $\frac{1}{m} + \frac{1}{n} < \frac{1}{2}$. In addition, the same result holds for $m = n = 4$.

(b). There is a unique rotationally symmetric graph-directed invariant forms on K_λ if $\frac{1}{m} + \frac{1}{n} > \frac{1}{2}$.

Remark. The above proposition follows directly from Sabot's Theorem once we have calculated the exact values of $\overline{\rho}, \underline{\rho}$'s for $\mathcal{J}_1$ and $\mathcal{J}_2$. The only unclear case is the critical case $\frac{1}{m} + \frac{1}{n} = \frac{1}{2}$. Indeed, there are 3 possible choices: $(m, n) = (3, 6), (6, 3)$ or $(4, 4)$. The $m = n = 4$ case can be studied directly by computation, which is tricky and long. We claim the non-existence result in this case without providing the details. For the $(m, n) = (3, 6)$ and $(6, 3)$ cases, experiments indicate that there is no solution to (4.3).

Example 4.2. *We can compute the unique rotationally invariant symmetric form on* K_λ *when* $m = 1$. *See Figure 14 for some typical such Julia sets. In particular, the Julia set corresponding to the* $m = 1, n = 2$ *case is homeomorphic to the double cover of the Sierpinski gasket, and as we shall see has the same renormalization constant.*



$m = 1, n = 2$ $m = 1, n = 3$

FIGURE 14. Some Julia sets with $m = 1$.

Computing the exact solution is tedious in general, but the renormalization constant η *is surprisingly concise:*

$$\eta = \frac{2n+1}{n+1}.$$

As we can see, for a rational map R_λ possessing a fixed critical point, its associated Julia set is quite different from those of MS maps. We leave the more general case, for example R_λ possessing a periodic critical point, for future studies.

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